

Wuhan Institute of Physics  
and Mathematics

April 25, 2017



# Soliton Collisions on the Edge of Integrability

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Boris Malomed (Tel Aviv)



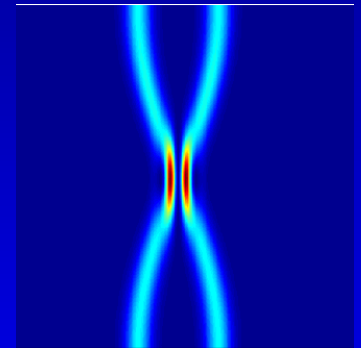
# Outline

- Intro to BEC with  $a < 0$

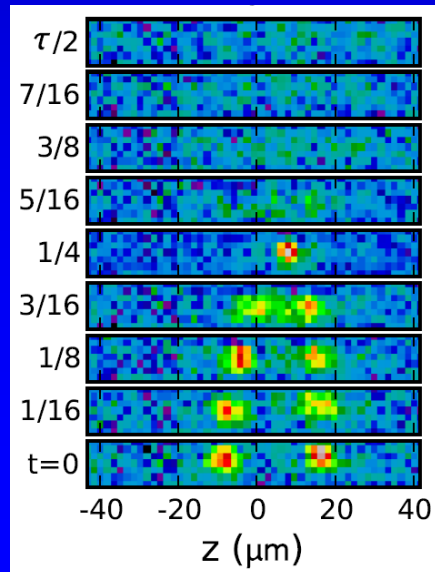


- Intro to solitons

- Phase-dependent collisions



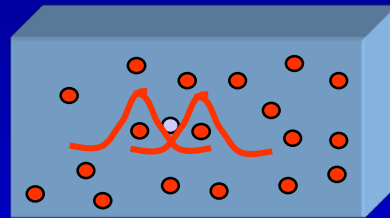
- Role of integrability



# Quantum Gases

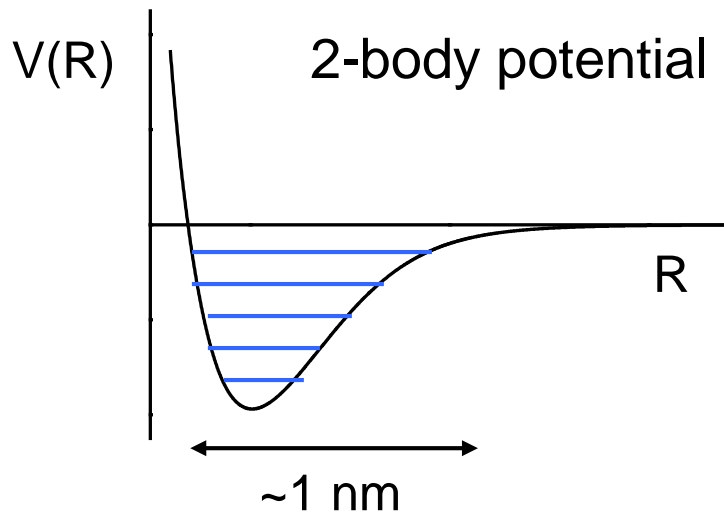
- Quantum regime  $n\Lambda^3 \geq 1$   $n$  = density  
 $\Lambda$  = de Broglie wavelength

Identical particles!



- Gas phase  $n \approx 10^{12} \text{ cm}^{-3}$
- Low temperature  $T \approx 100 \text{ nK}$   
 $\Rightarrow \Lambda \approx 1 \text{ }\mu\text{m}$
- Phase transitions
  - Bosons ( $^7\text{Li}$ ): Bose-Einstein condensation
  - Fermions ( $^6\text{Li}$ ): Fermion pairing

# Interactions - Generic Discussion



$$\Lambda_{\text{dB}} \sim 1 \text{ } \mu\text{m}$$

$\Rightarrow$  detailed shape of  $V(R)$   
unimportant

Characterize interaction by s-wave scattering length  $a$ :

- mean-field interaction energy  $nU_0 = 4\pi\hbar^2 n a / m$

$a < 0$  attractive

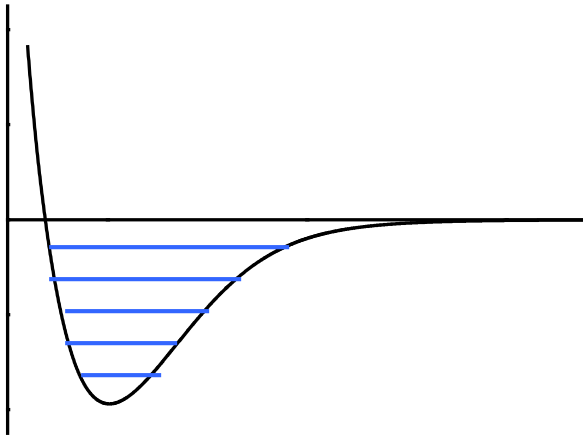
$a > 0$  repulsive

# Implications of Interactions

- Bosons
  - stability of condensate
  - bright or dark solitons
  - healing length: vortices, speed of sound
  - excitation spectrum
  - Mott insulator: on-site interactions  $U$
  - miscibility or immiscibility of spinor condensates
- Fermions
  - Cooper pairing (BCS) or molecules (BEC)
  - Hubbard model:  $U/t$

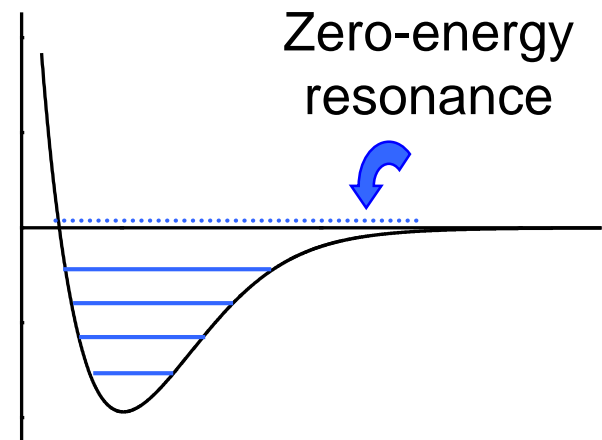
# What Determines $a$ ?

${}^7\text{Li}$



$$a = -27 a_0$$

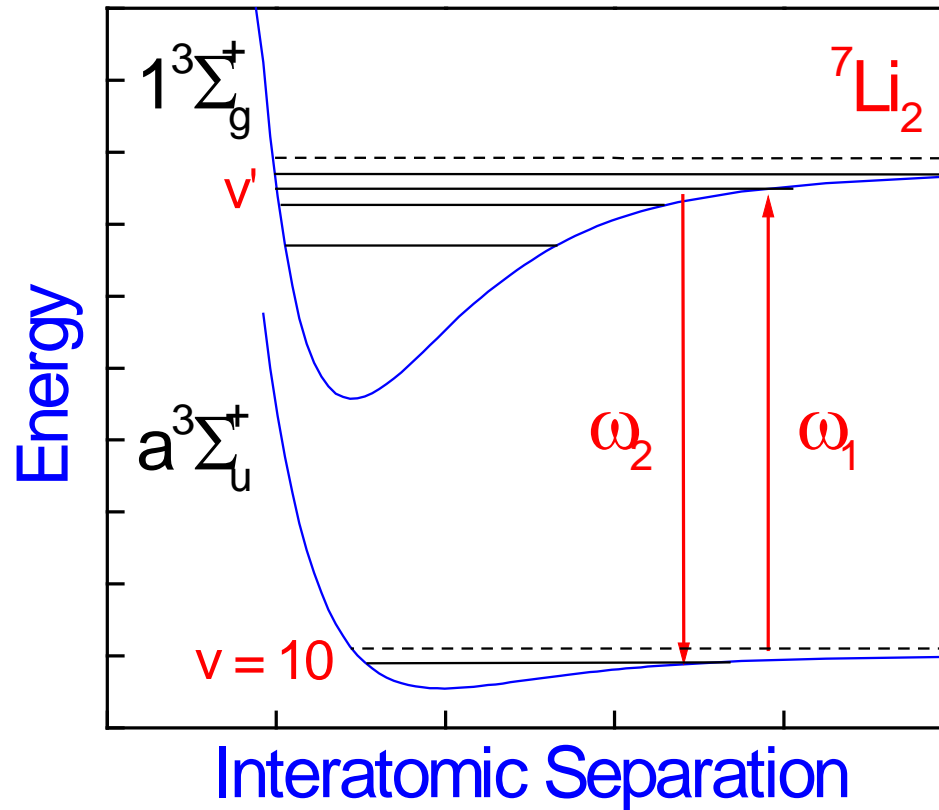
${}^6\text{Li}$



$$a = -2300 a_0$$

Answer: The last bound state!

# Measuring $a$ by Photoassociation



Abraham *et al.*,  
PRL **74**, 1315 (1995)

2-photon photoassociation:  $E_B = \omega_2 - \omega_1$

Li potentials well characterized: 2-body physics known *precisely*

# Implications of Interactions for BEC

Mean-field interaction energy  $U = 4\pi\hbar^2 a n / m$

$a$  is the s-wave scattering length

$a > 0$  (eg H,  $^{23}\text{Na}$ ,  $^{87}\text{Rb}$ ): **repulsive**

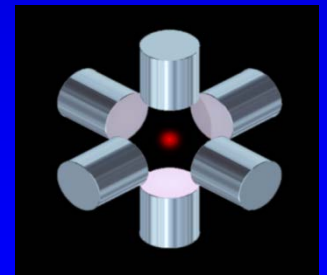
- Stable BEC

$a < 0$  (eg  $^7\text{Li}$ ,  $^{85}\text{Rb}$ ): **attractive**

- $dU/dn < 0 \Rightarrow$  mechanically unstable

$\therefore$  BEC in a gas thought not possible

(Bogolubov, 1947; Landau and Lifshitz, 1958;  
Stoof, 1994)



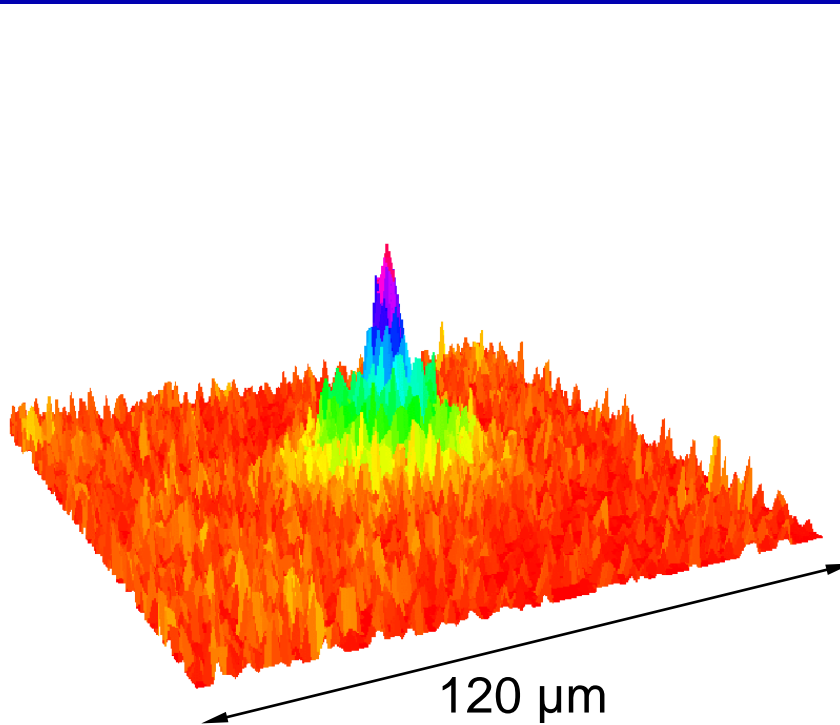
Permanent Magnet  
Trap 1995

**Attractive condensate predicted to implode**

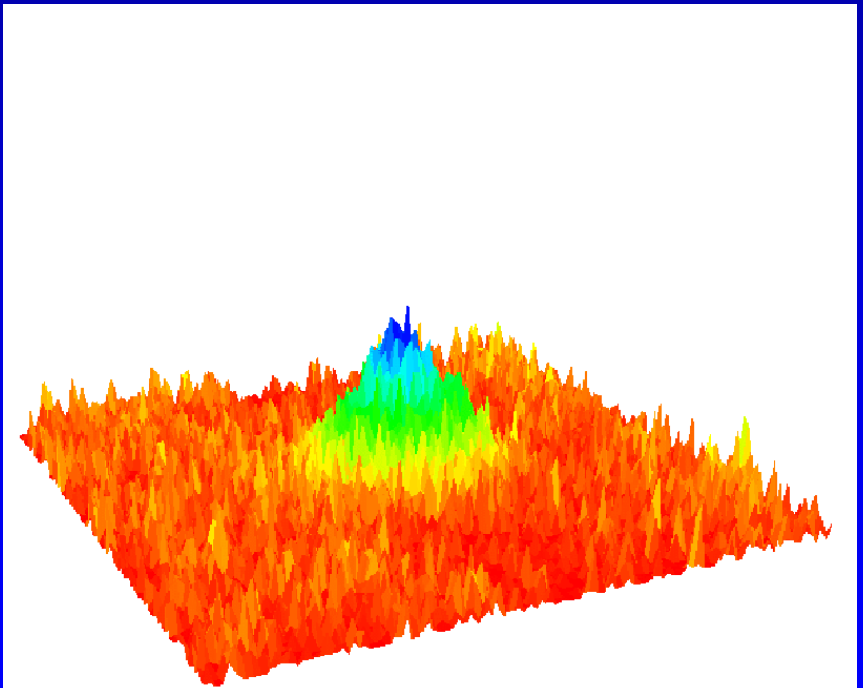


# Result

There are condensates!  
(although they are really puny)



$N = 23\text{k}, N_0 = 1050$

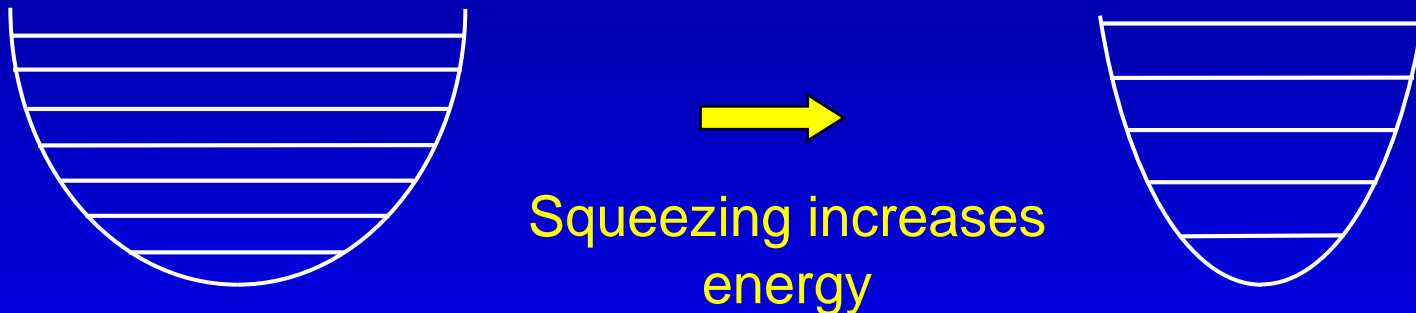


$N = 23\text{k}, N_0 = 65$

# BEC with Attractive Interactions - ~~3D~~ 0D

Mean-field interaction energy  $U = 4\pi\hbar^2 a n / m$

- Condensate mechanically unstable – no BEC in free space
- Stabilized by quantum pressure in a trap

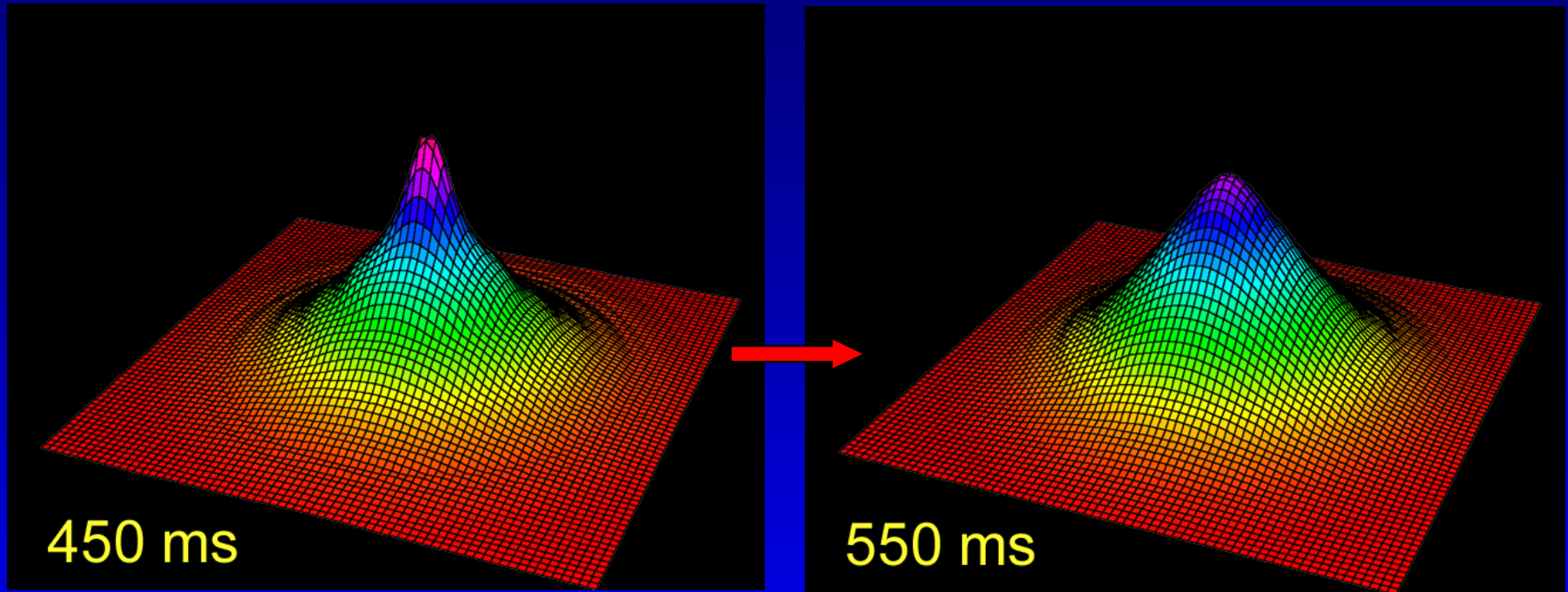


- Attraction balanced by zero-point energy

BEC is possible with  $a < 0$ , but  $N_0$  limited:

$$U < \hbar\omega_{r,z} \longrightarrow \text{"0D" limit}$$

# Quench Cool and Attractive Condensate



Collapse!

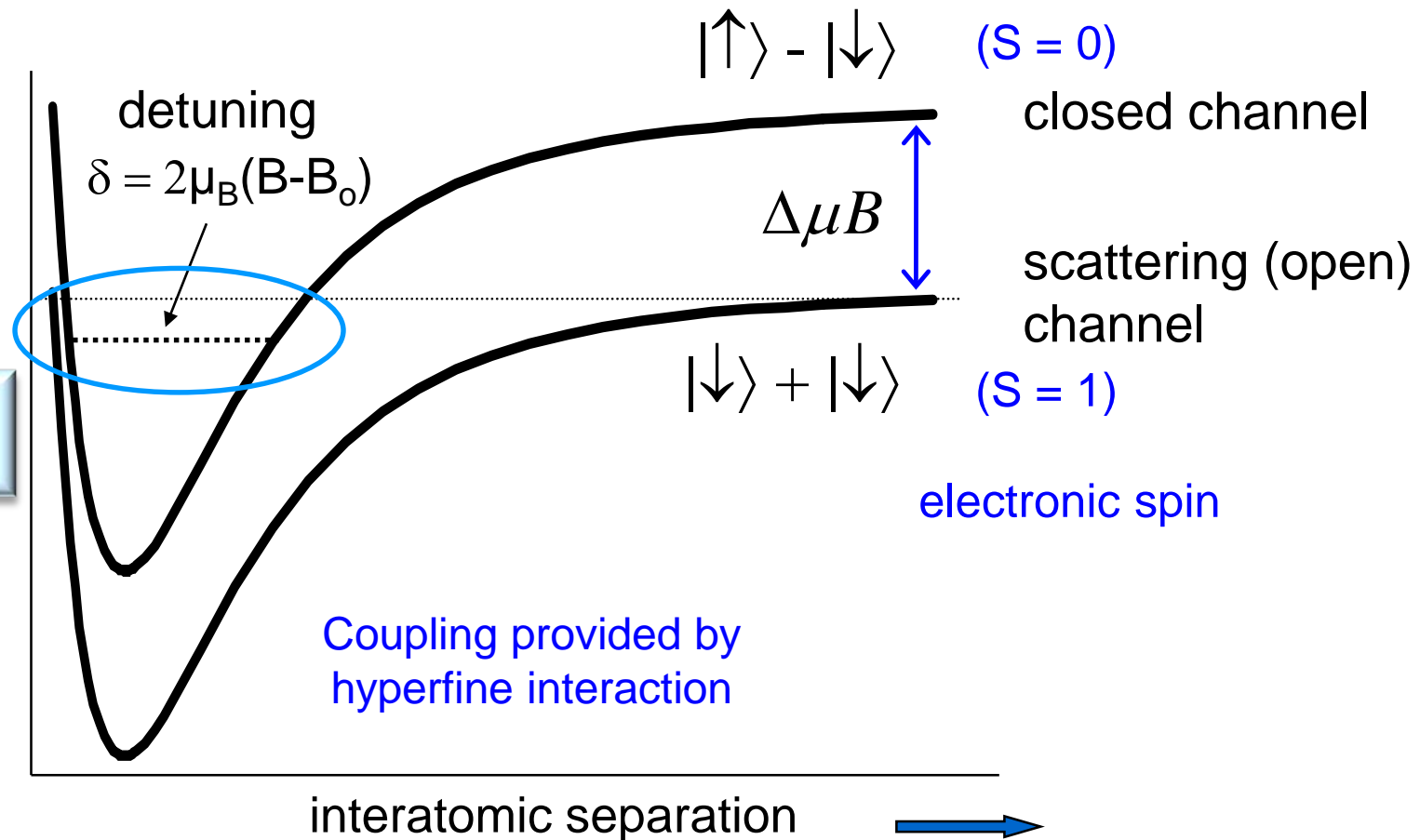
*“Bosenova”*

Gerton *et al.*, *Nature* **408**, 692 (2000) - Rice  
Donley *et al.*, *Nature* **412**, 295 (2001) - JILA

# Tunable Interactions - Feshbach Resonance

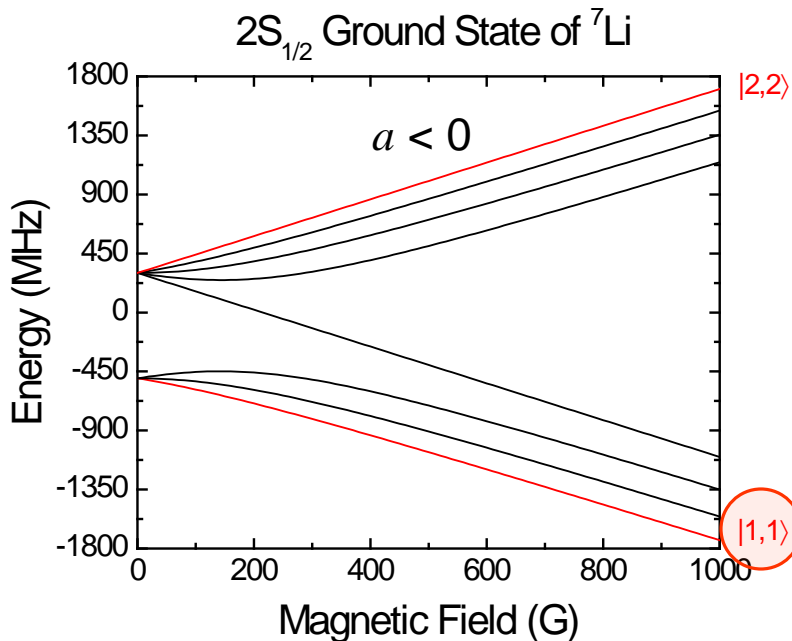
Magnetically tune free atoms into resonance with a bound molecular state:

Alkali metal atoms interact via singlet or triplet potential

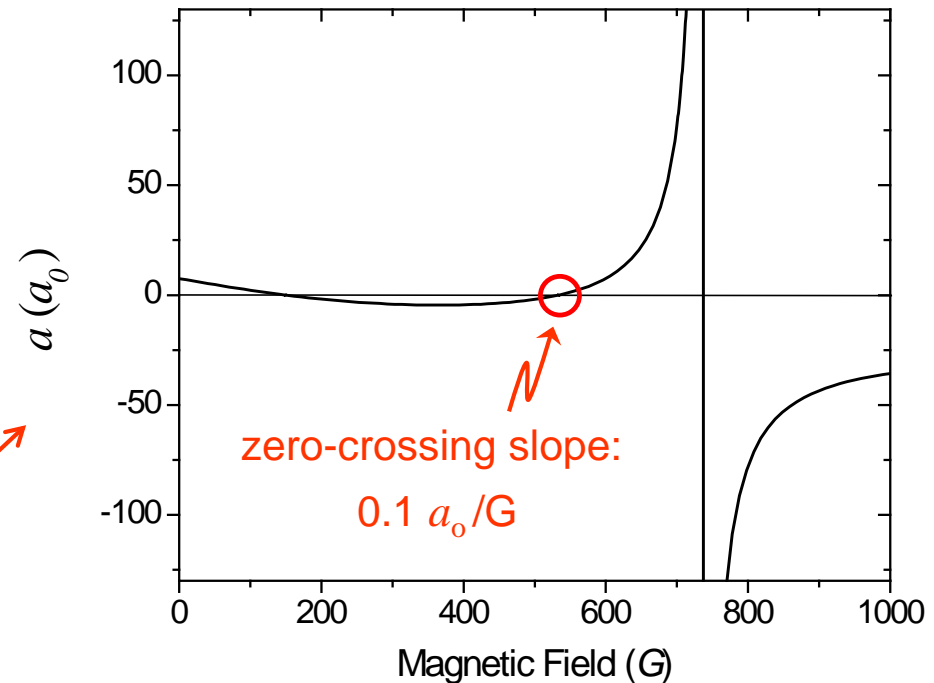


# Feshbach Resonance in $^7\text{Li}$ (Boson)

Hyperfine sublevels of  $^7\text{Li}$



Coupled channels calculation  
of the scattering length of  
 $^7\text{Li}$   $|1,1\rangle$  state



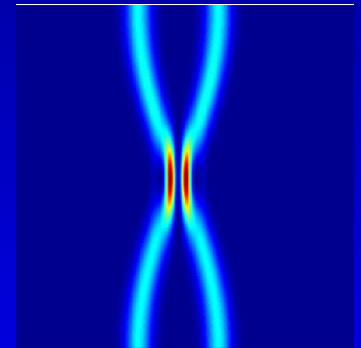
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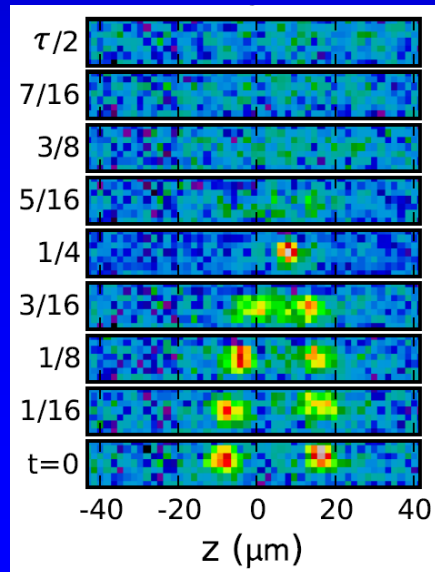


- Intro to solitons

- Phase-dependent collisions



- Role of integrability



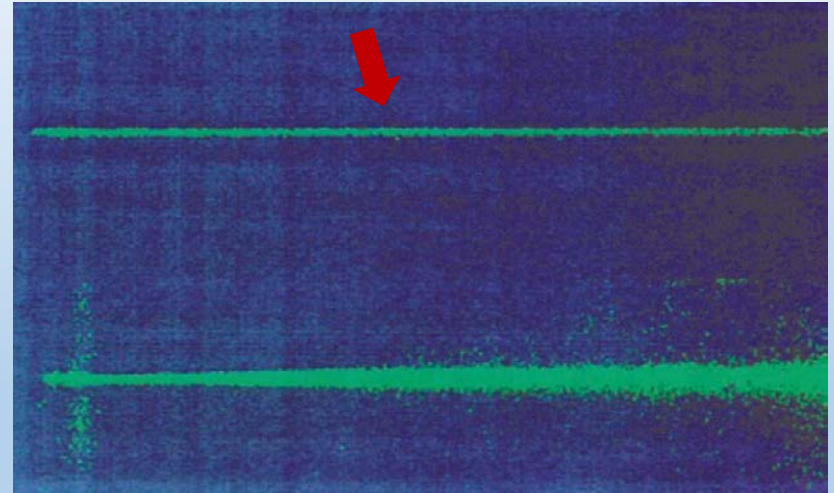


# Solitons are Everywhere!

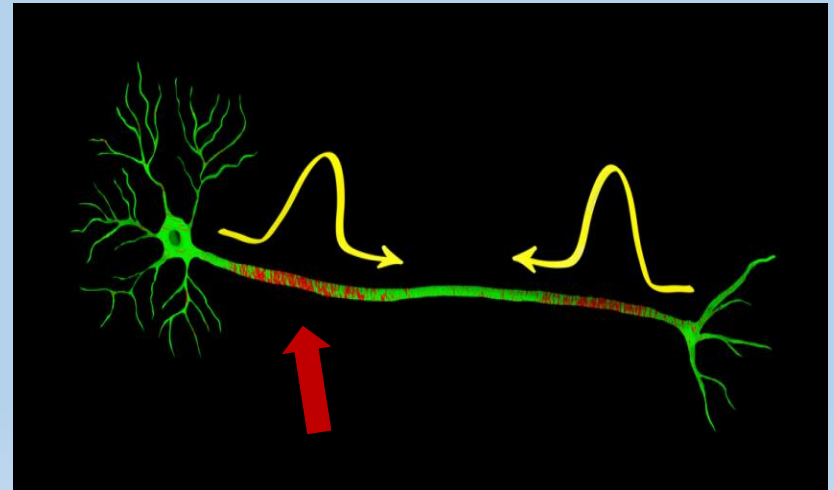
Scott Russell Aqueduct - Edinburgh



Photorefractive Crystal



Gulf of Carpentaria, Australia



Nerve Impulses – soliton collisions?

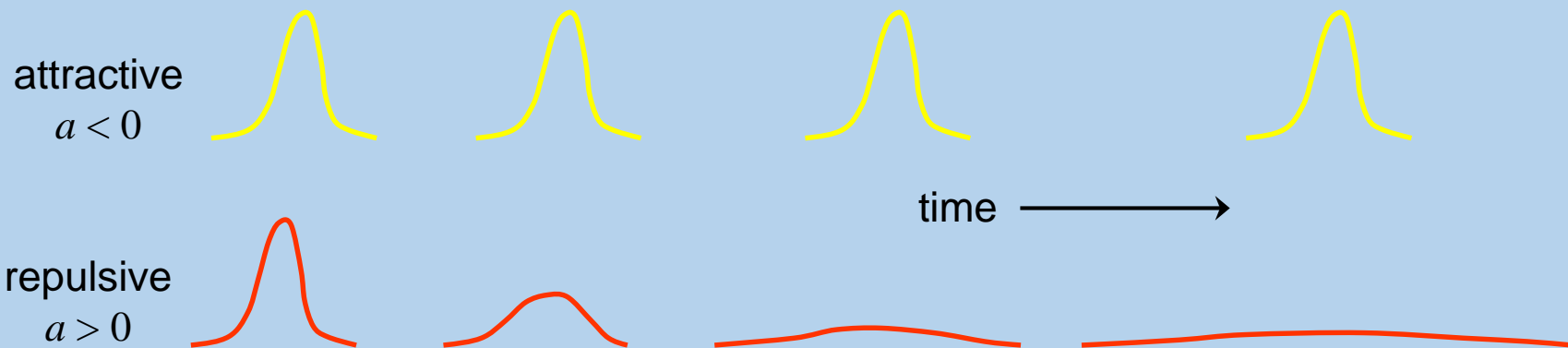
# General Properties of Solitons

- mathematically described by 1D nonlinear partial differential equations that are integrable, e.g

$$i\hbar \frac{d}{dt} \Psi = \left( -\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + g|\Psi|^2 \right) \Psi$$

where  $g = 2\hbar^2/ma_{1D} < 0$   
and  $a_{1D} = a_r/a$

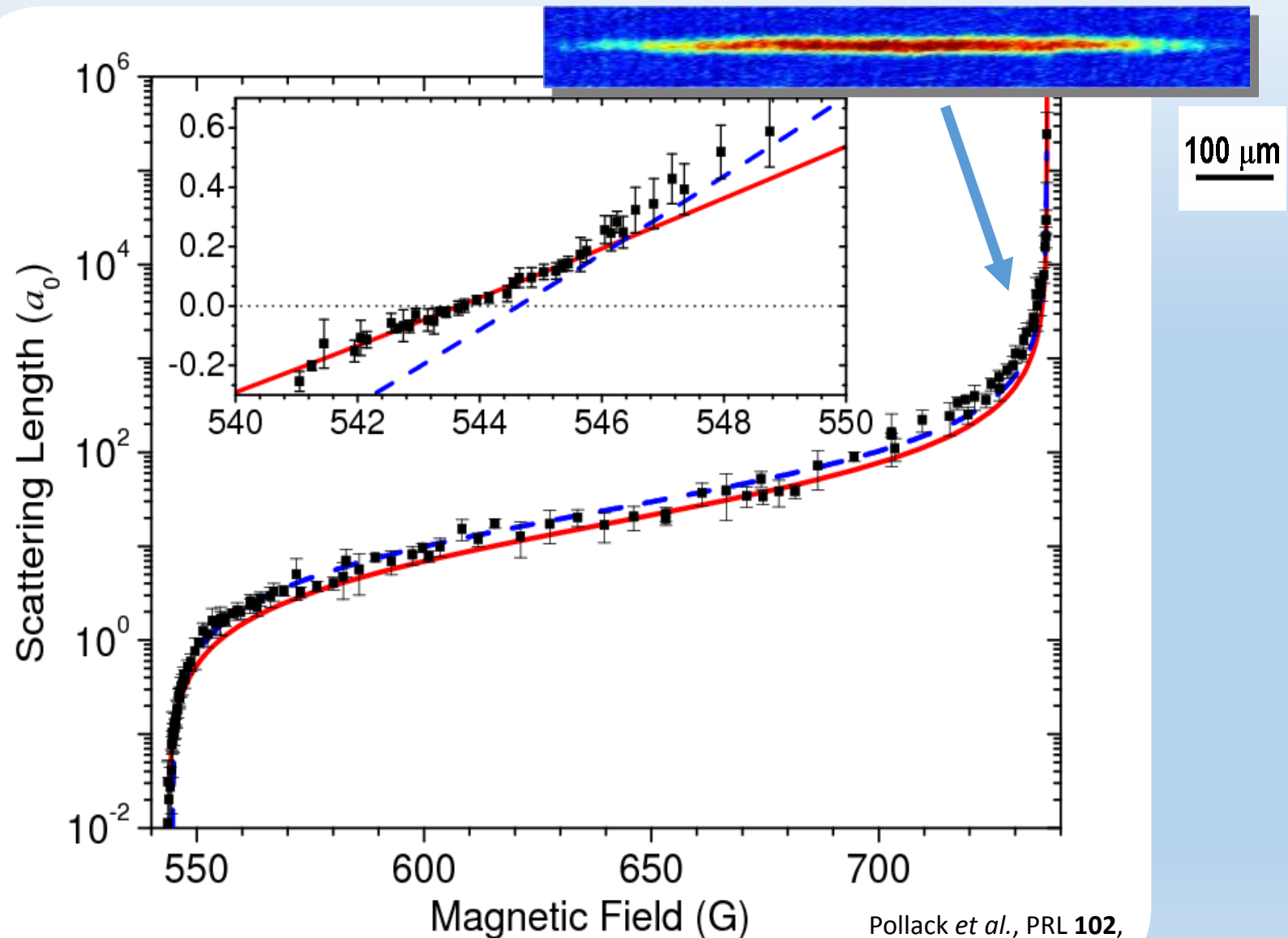
- dispersion compensated by nonlinear interaction – no spreading



- survives collisions without change in shape, amplitude or velocity (except for a possible phase jump)
  - independent of the number of collisions
  - system does not *thermalize*

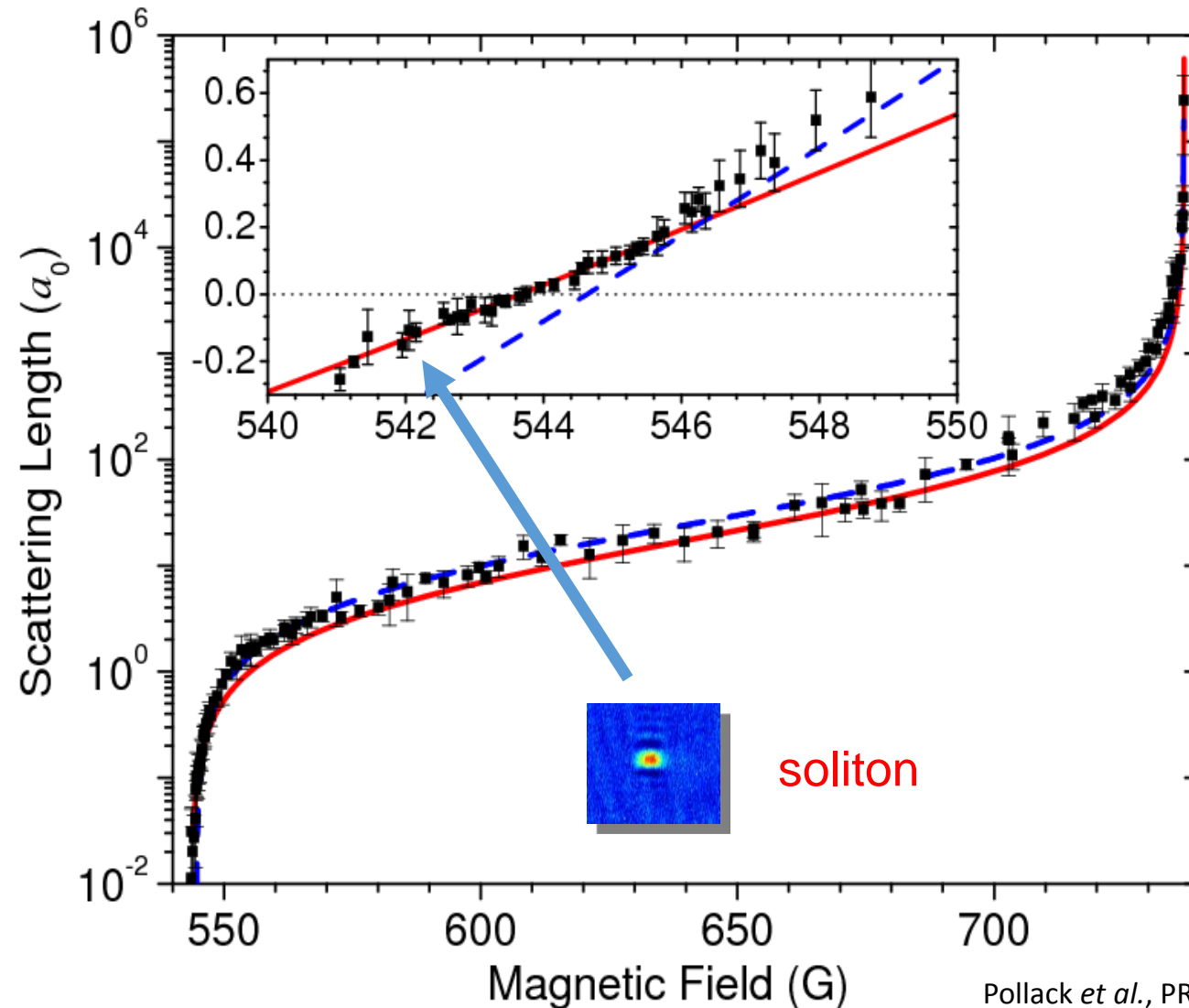


# Feshbach Resonance in $^7\text{Li}$ – Stable BEC



Pollack *et al.*, PRL **102**,  
090402 (2009)

# Sweep Through the Zero-Crossing to Make Soliton(s)



Pollack *et al.*, PRL **102**,  
090402 (2009)

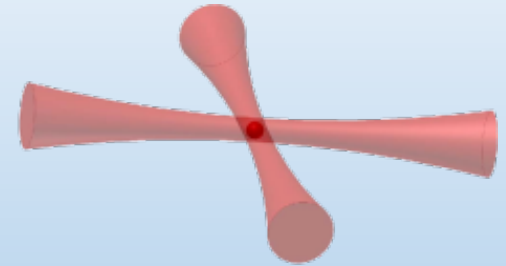
c.f. Khaykovich *et al.*, Science **296**, 1290 (2002) *ENS*;  
Strecker *et al.*, Nature **417**, 150 (2002) *Rice*

# Expansion into 1D Waveguide – Nondispersive

Formed in crossed beam trap and transferred into single beam:

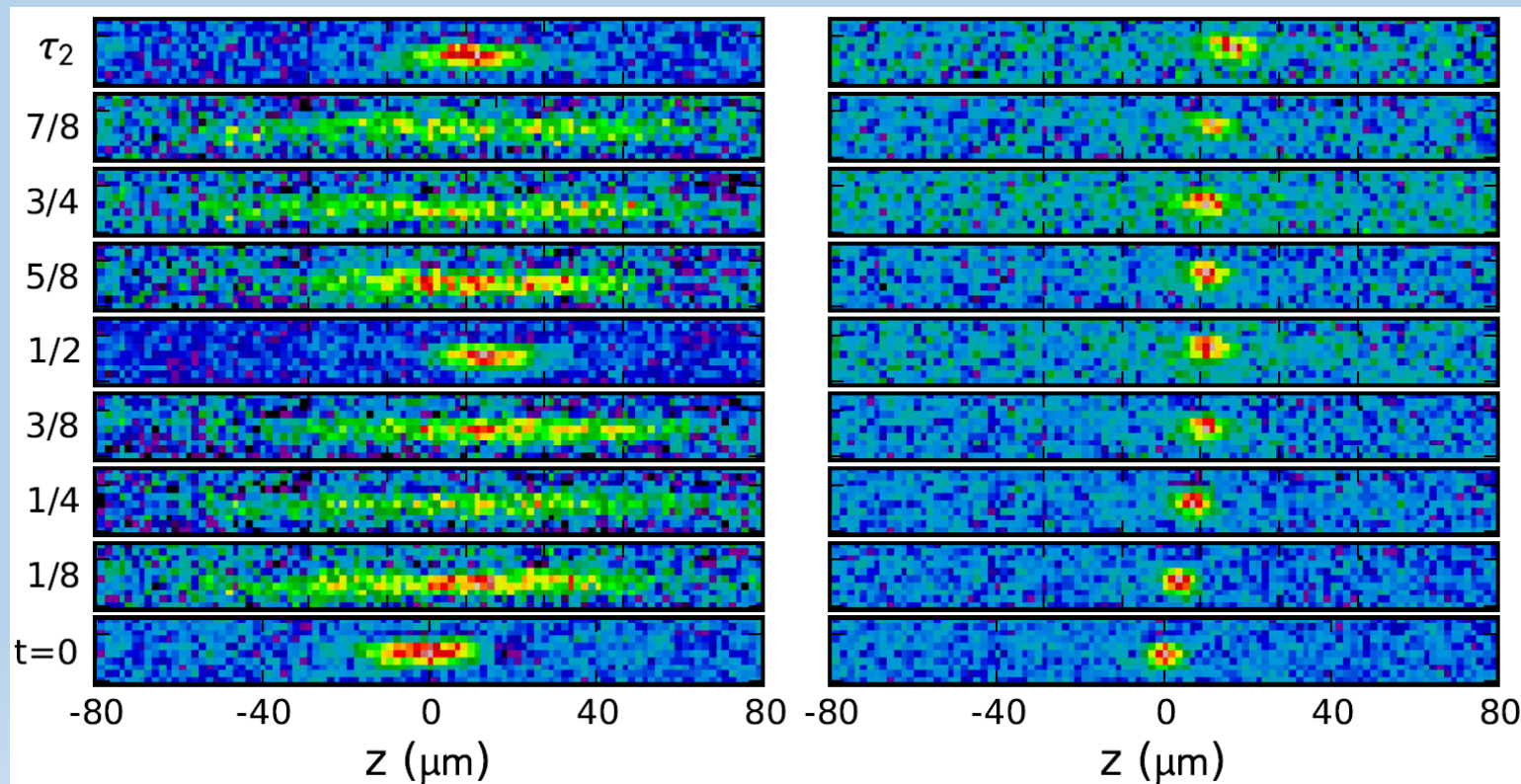
Crossed beam:  $\omega_z = 2\pi \times 31 \text{ Hz}$  ( $\tau = 32 \text{ ms}$ )

Single beam:  $\omega_z = 2\pi \times 8 \text{ Hz}$  ( $\tau = 125 \text{ ms}$ )



$N/N_c = +0.5$

$N/N_c = -0.5$



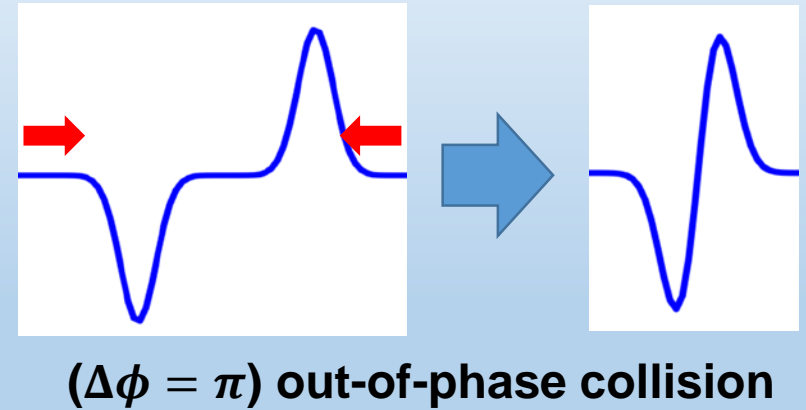
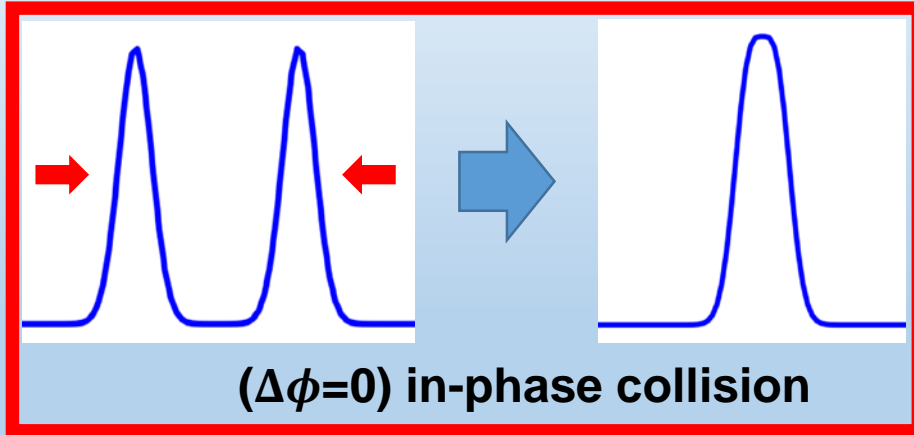
Stable only in  
quasi-1D  $\rightarrow$   
critical number:

$$N_c = 0.67 \frac{a_r}{|a|}$$

$$a_r = \sqrt{\frac{\hbar}{m \omega_r}}$$

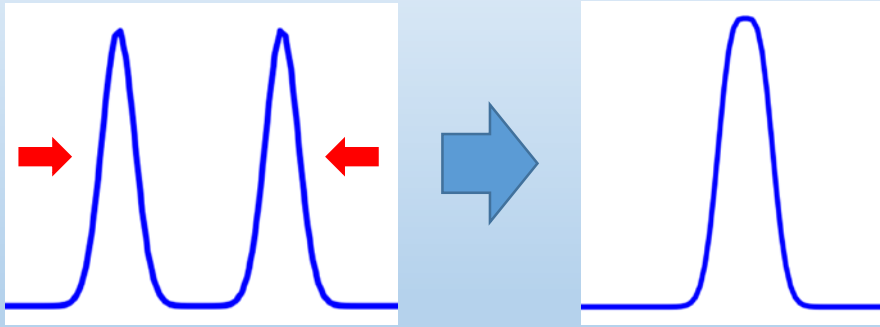
# Phase-dependent Interactions

- **Gordon-Haus Effect:** J. P. Gordon, Opt. Lett. 8, 596 (1983)

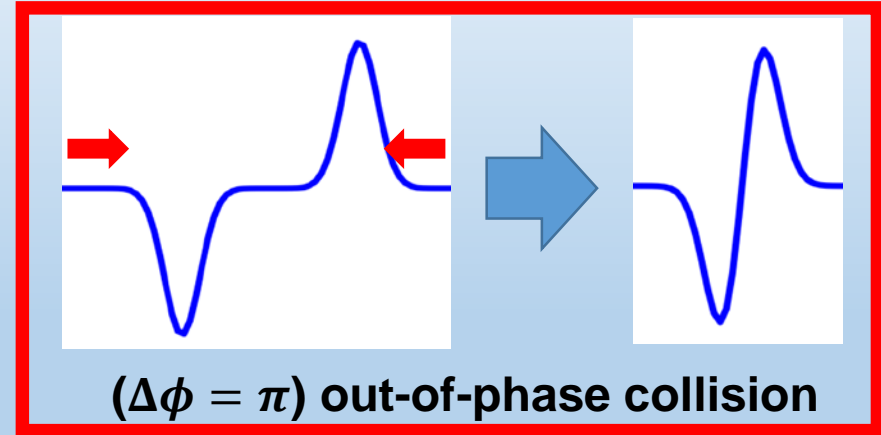


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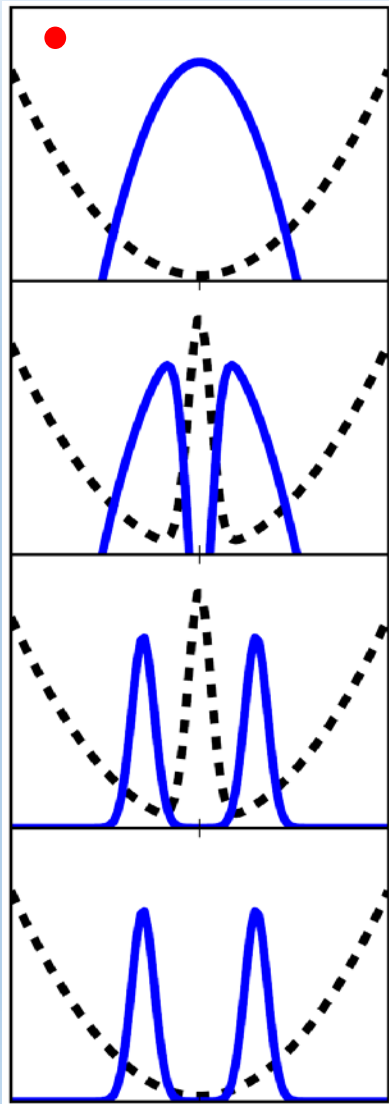


$(\Delta\phi=0)$  in-phase collision

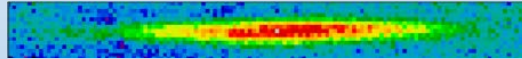


$(\Delta\phi = \pi)$  out-of-phase collision

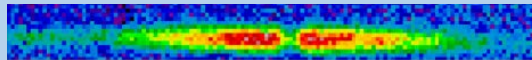
# Formation of Soliton Pair for Collision Study



- form BEC by evaporation at  $+140a_0$



- turn on barrier



- ramp magnetic field to  $-0.57a_0$

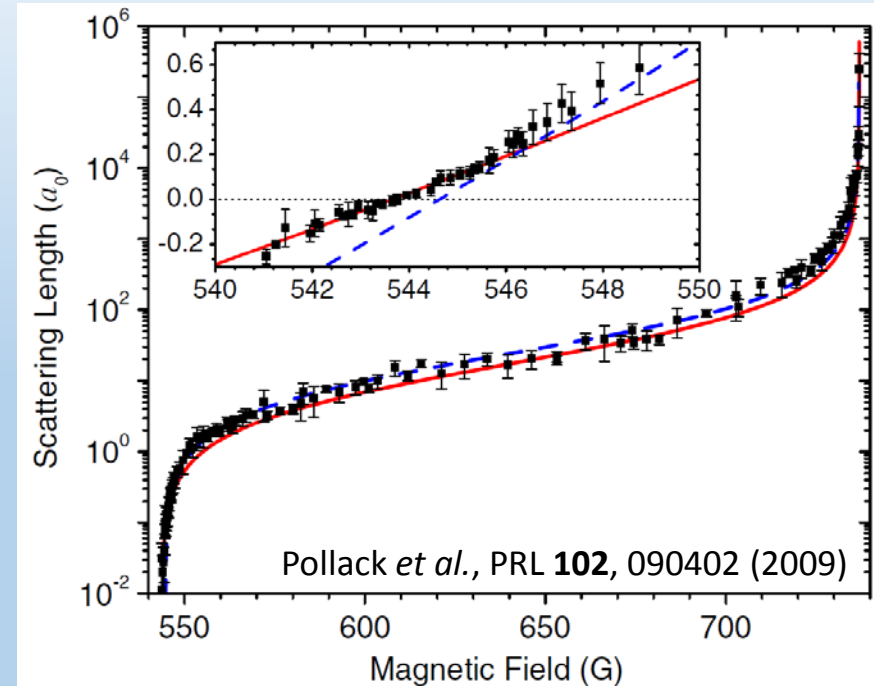


- quickly turn barrier off



$$\Delta z = 26 \mu m$$

$^7\text{Li}$  Feshbach Resonance

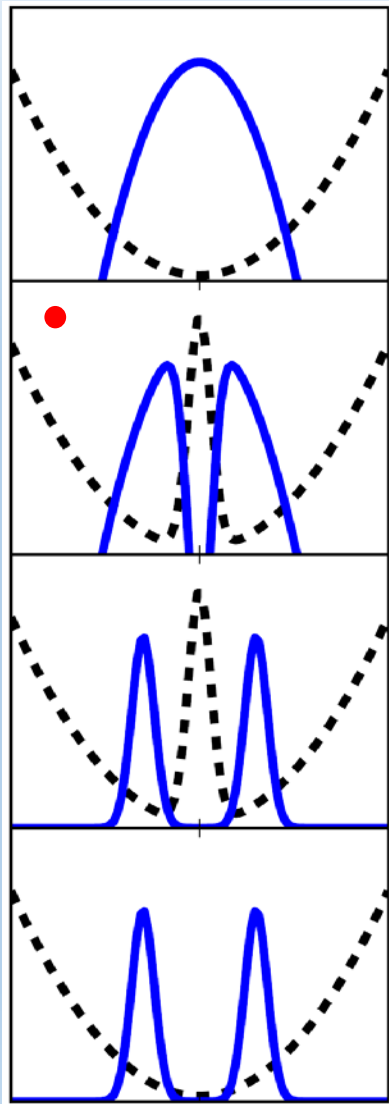


Atoms in  $|F = 1, m_F = 1\rangle$  state

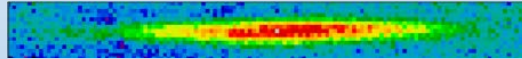
$$\omega_z = 2\pi \times 31 \text{ Hz}$$

$$\omega_r = 2\pi \times 254 \text{ Hz}$$

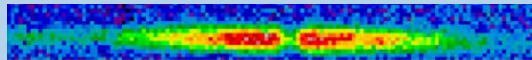
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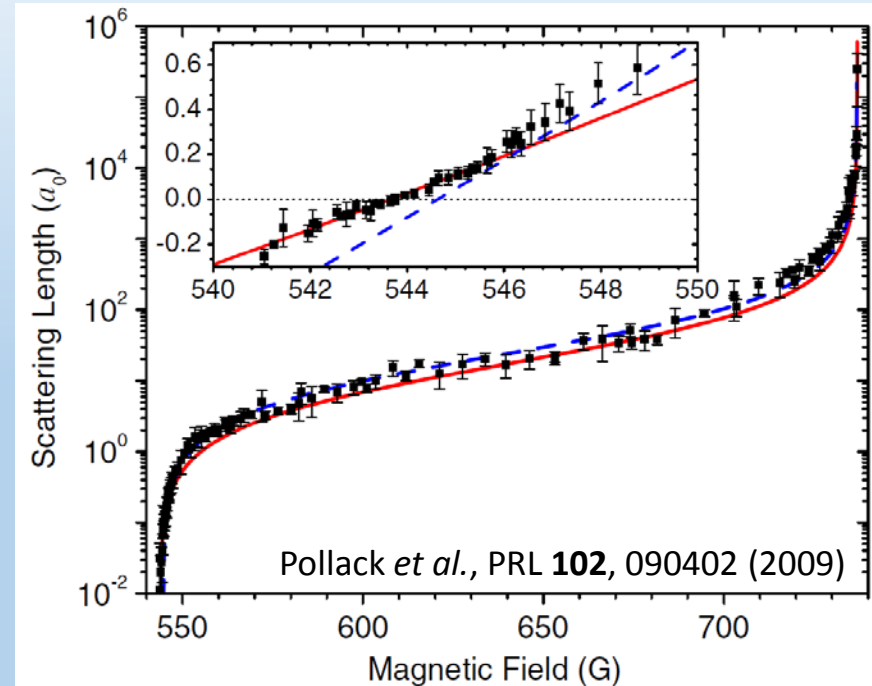
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$^7\text{Li}$  Feshbach Resonance



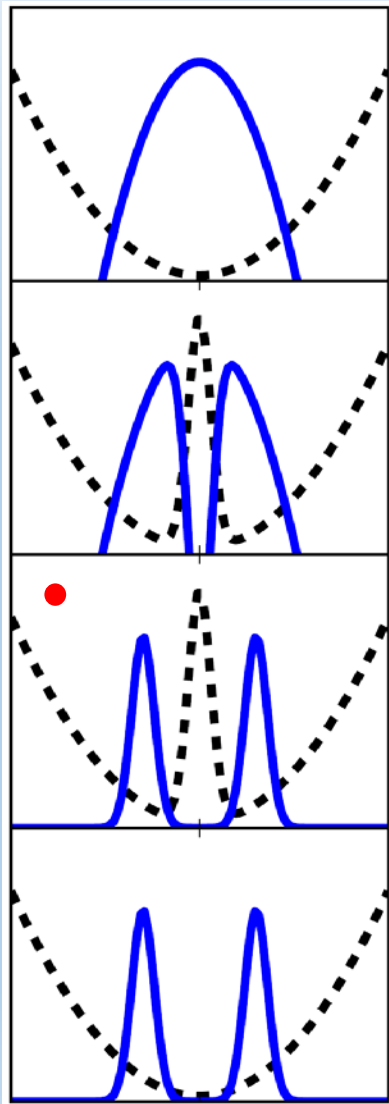
Atoms in  $|F = 1, m_F = 1\rangle$  state

$$w_x = 2.2 \text{ mm}$$

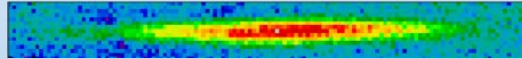
$$w_z = 5.6 \mu\text{m}$$

900 GHz blue-detuned light sheet

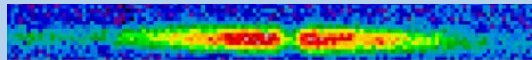
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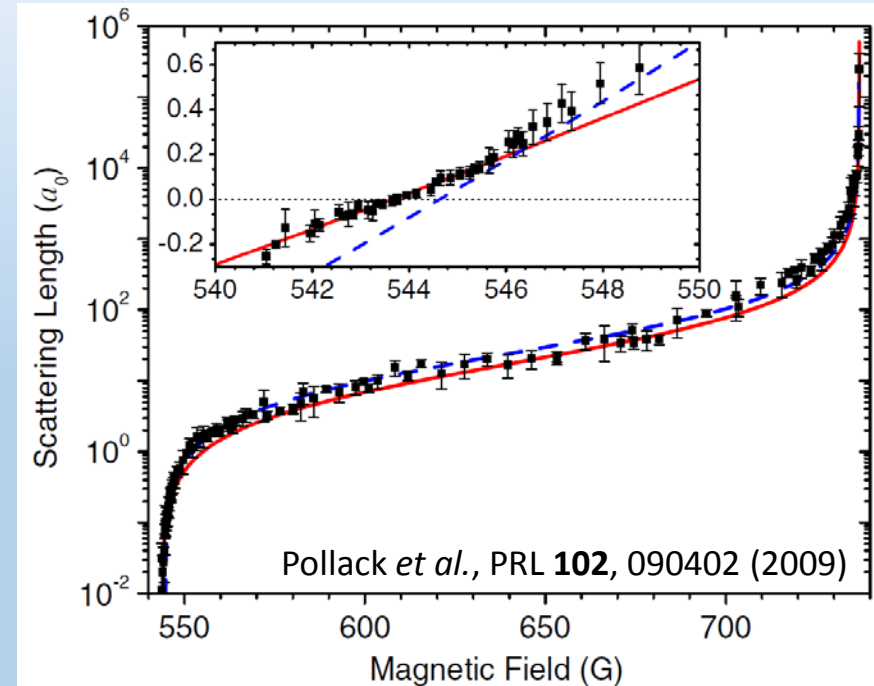
- ramp magnetic field to  $-0.57a_0$



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$^7\text{Li}$  Feshbach Resonance



Atoms in  $|F = 1, m_F = 1\rangle$  state

$$\omega_z = 2\pi \times 31 \text{ Hz}$$

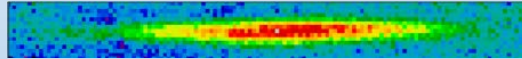
$$\omega_r = 2\pi \times 254 \text{ Hz}$$

$$N \sim 28\,000, N/N_c = -0.53$$

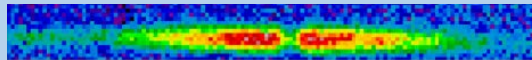


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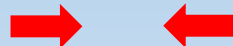
- turn on barrier



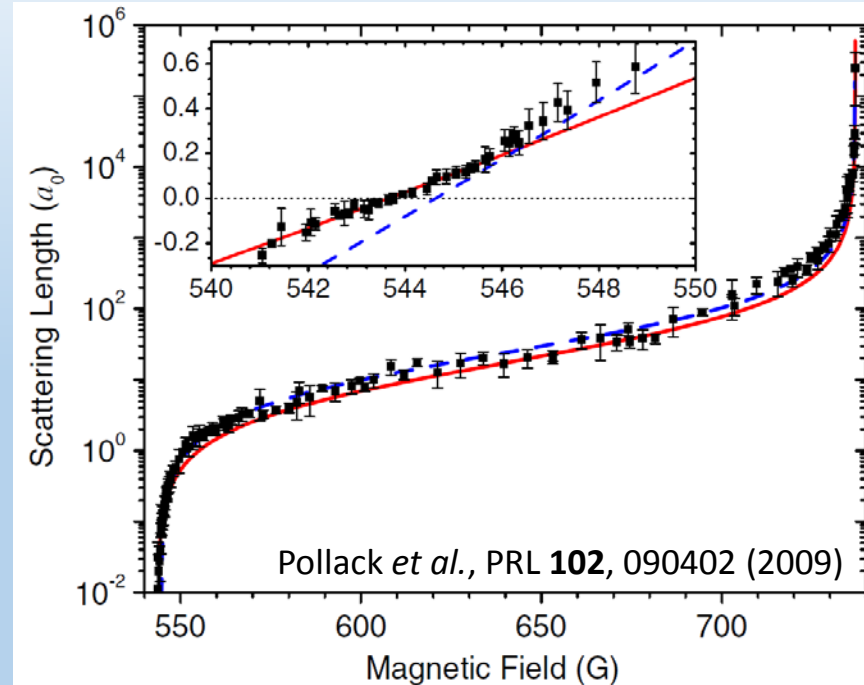
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- quickly turn barrier off



$^7\text{Li}$  Feshbach Resonance



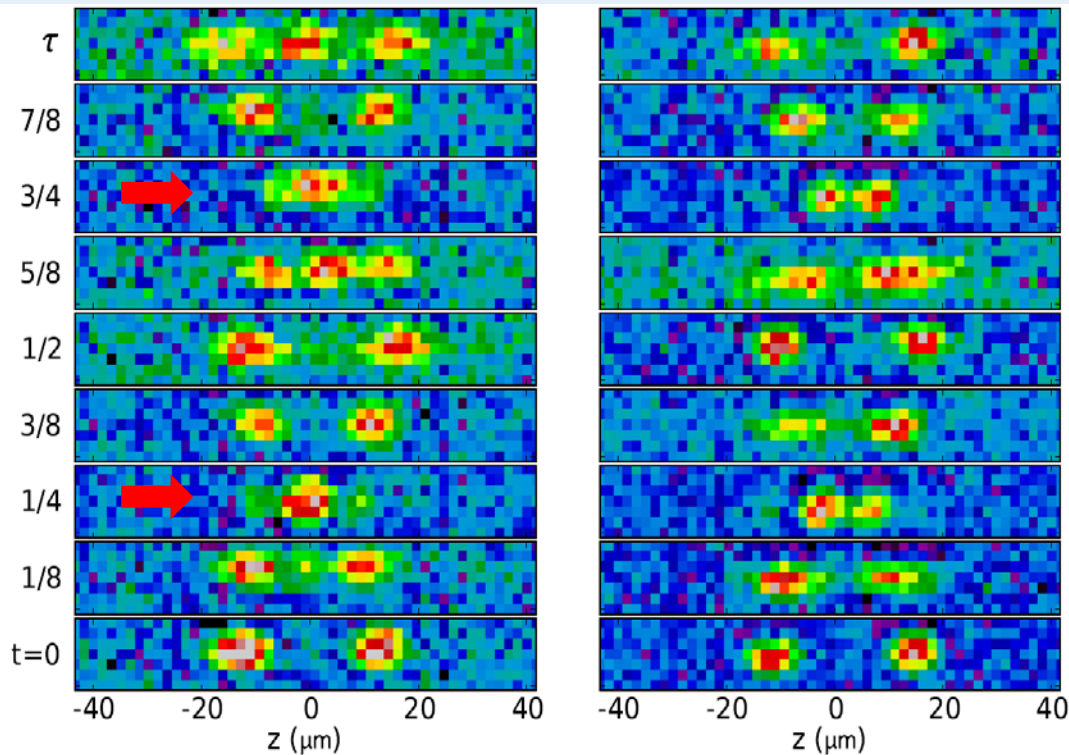
Atoms in  $|F = 1, m_F = 1\rangle$  state

$$\omega_z = 2\pi \times 31 \text{ Hz}$$

$$\omega_r = 2\pi \times 254 \text{ Hz}$$

$$N \sim 28\,000, N/N_c = -0.53$$

# Phase-Dependent Collisions



$$\Delta\phi \approx 0$$

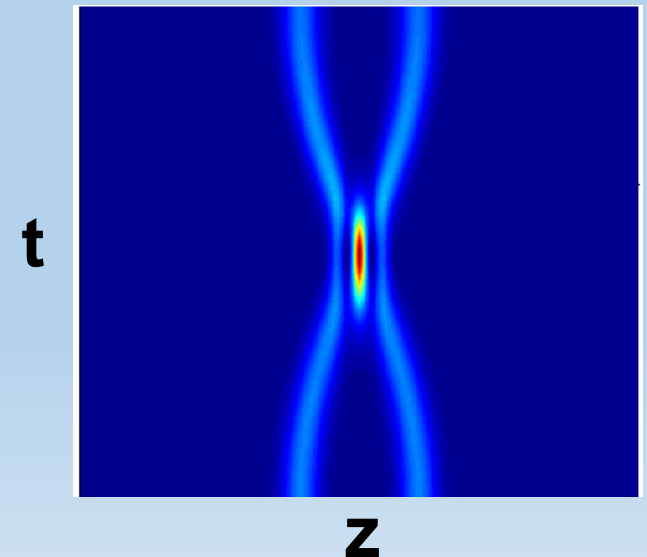
$$\Delta\phi \approx \pi$$

Nguyen, Dyke, Luo, Malomed & Hulet., Nat. Phys. **10**, 918-922 (2014)

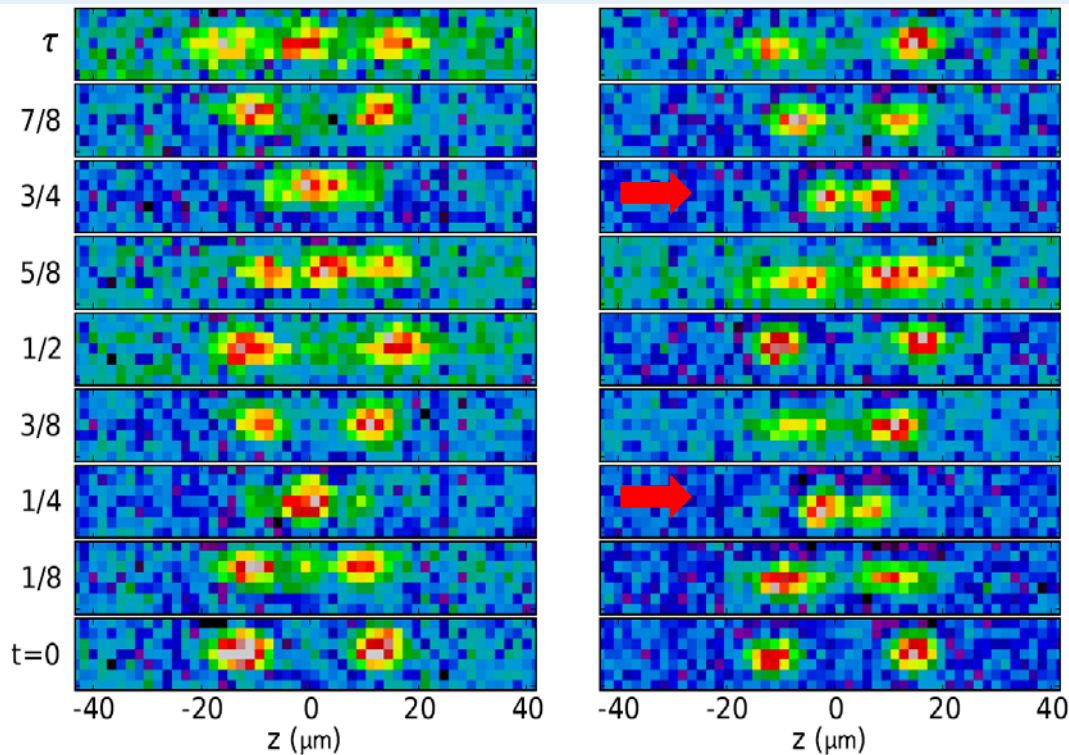
Also, Parker, Martin, Cornish, Adams,  
J Phys B 41, 045303 (2008)

- collisions for a full trap period ( $\tau = 32 \text{ ms}$ )
- multiple images using phase contrast
- $N/N_c = -0.53$
- $\Delta\phi$  inferred from simulations:

$\Delta\phi = 0$  collision



# Phase-Dependent Collisions



$$\Delta\phi \approx 0$$

$$\Delta\phi \approx \pi$$

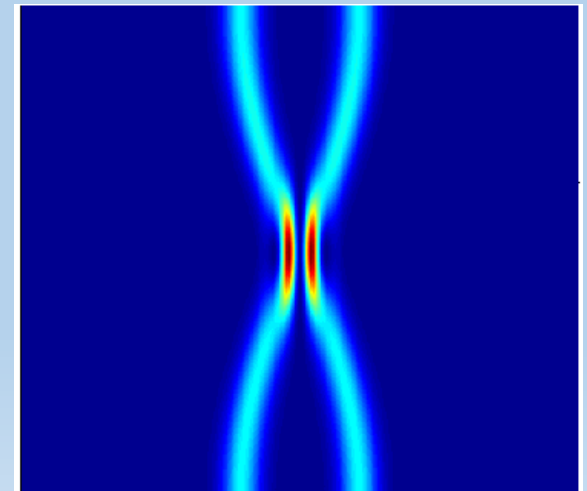
Nguyen, Dyke, Luo, Malomed & Hulet., Nat. Phys. **10**, 918-922 (2014)

But, integrability  $\rightarrow$   
solitons must pass  
through one another (?)

- collisions for a full trap period ( $\tau = 32 \text{ ms}$ )
- multiple images using phase contrast
- $N/N_c = -0.53$
- $\Delta\phi$  inferred from simulations:

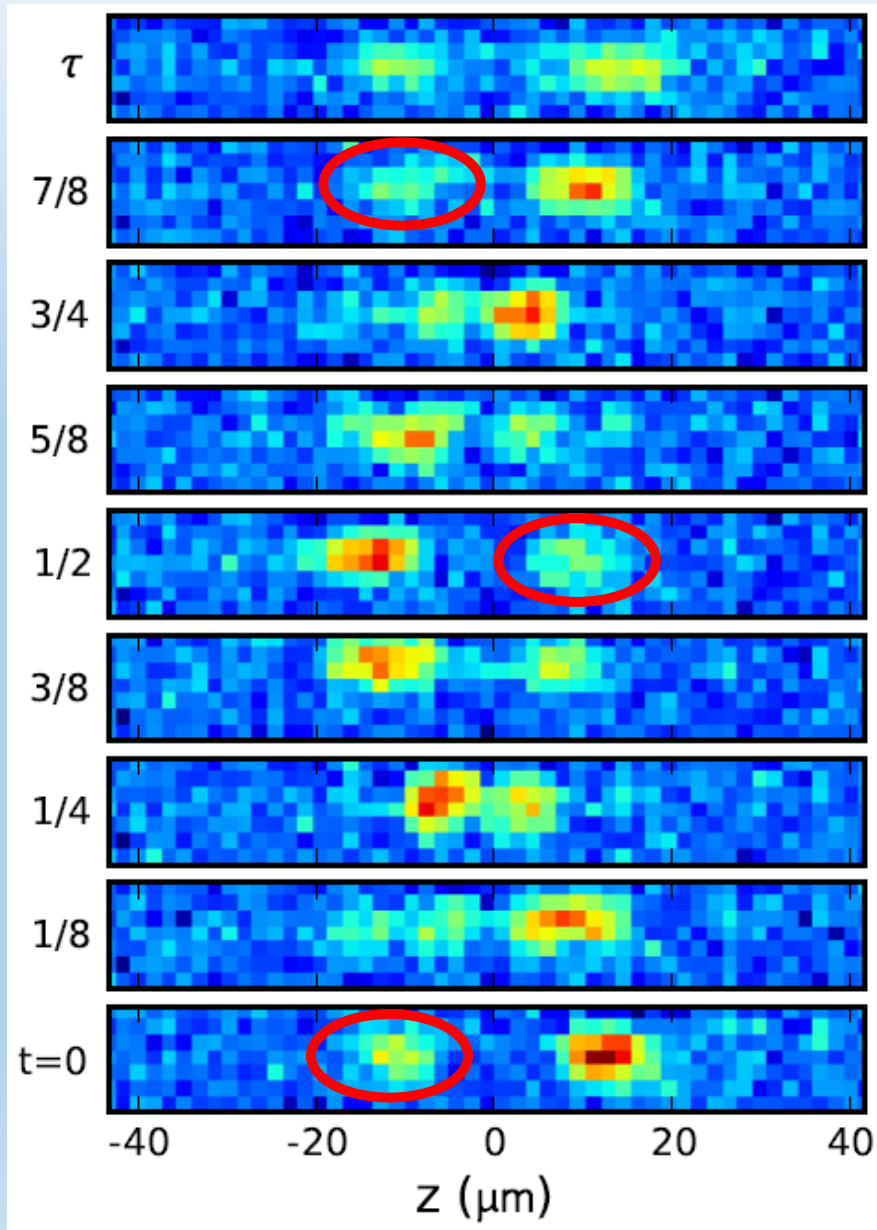
$$\Delta\phi = \pi \text{ collision}$$

**t**

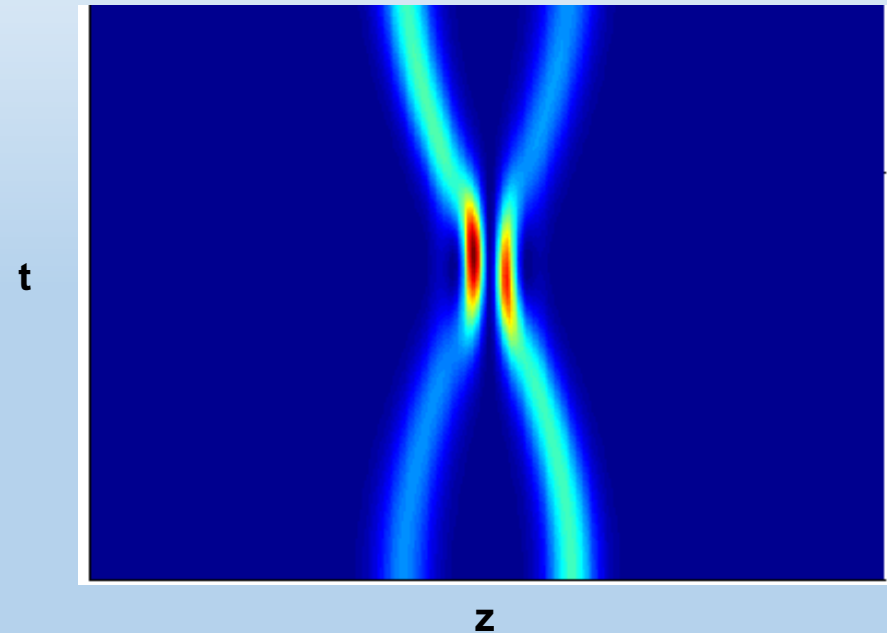


**z**

# Do They Cross? Tagged Collisions



- numerical simulation (2:1 ratio):



- use resonant beam to remove atoms from only one side
  - $\Delta\phi = \pi$ : appear to repel
  - solitons pass through one another

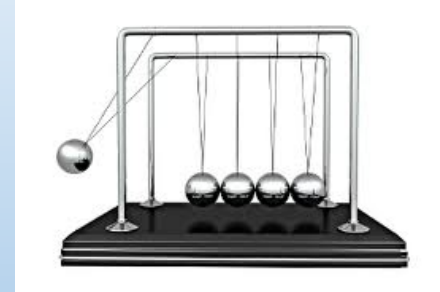
Consistent with integrability!

# Examples of Quantum Integrability

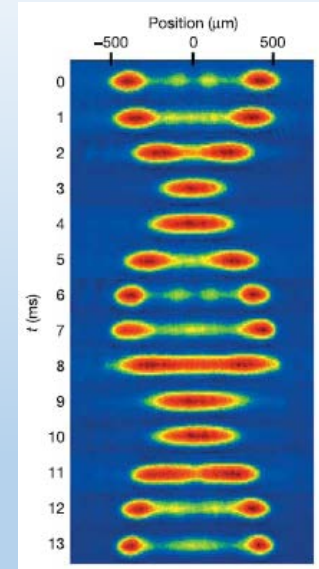
- **Lieb-Liniger Model:** 1D bosons with point interactions

“Quantum Newton’s Cradle”:

Kinoshita, Wenger, Weiss, Nature (2006)



100's of collisions  
without thermalization



- **1D solitons:**  $g|\psi|^3$  nonlinearity with  $g < 0$ , e.g. BEC with  $a < 0$

Quasi-1D if  $\mu > \frac{1}{2}\hbar\omega_r$ , otherwise unstable to *collapse*

The critical number for collapse is:  $N_c = 0.7 a_r / a$ , where  $a_r = (\hbar m \omega_r)^{1/2}$

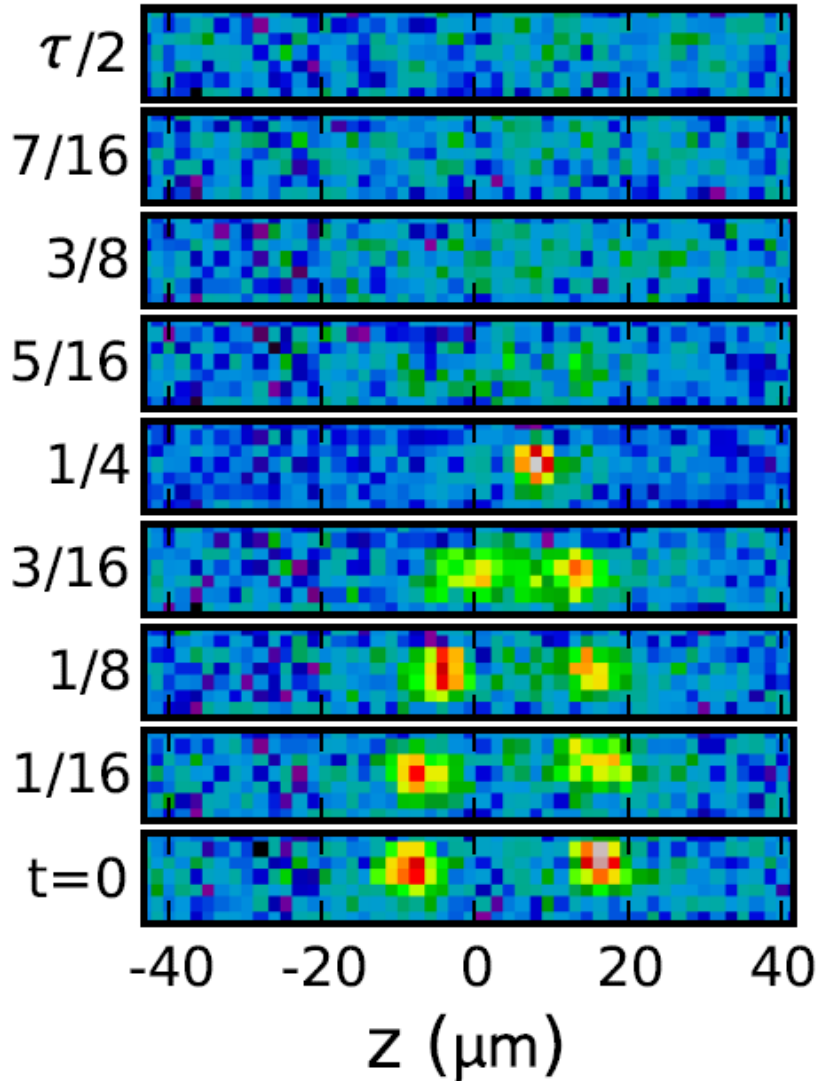
$N / N_c$  is a measure of the strength of the nonlinearity and

$N / N_c = 1$  defines an *integrability edge*

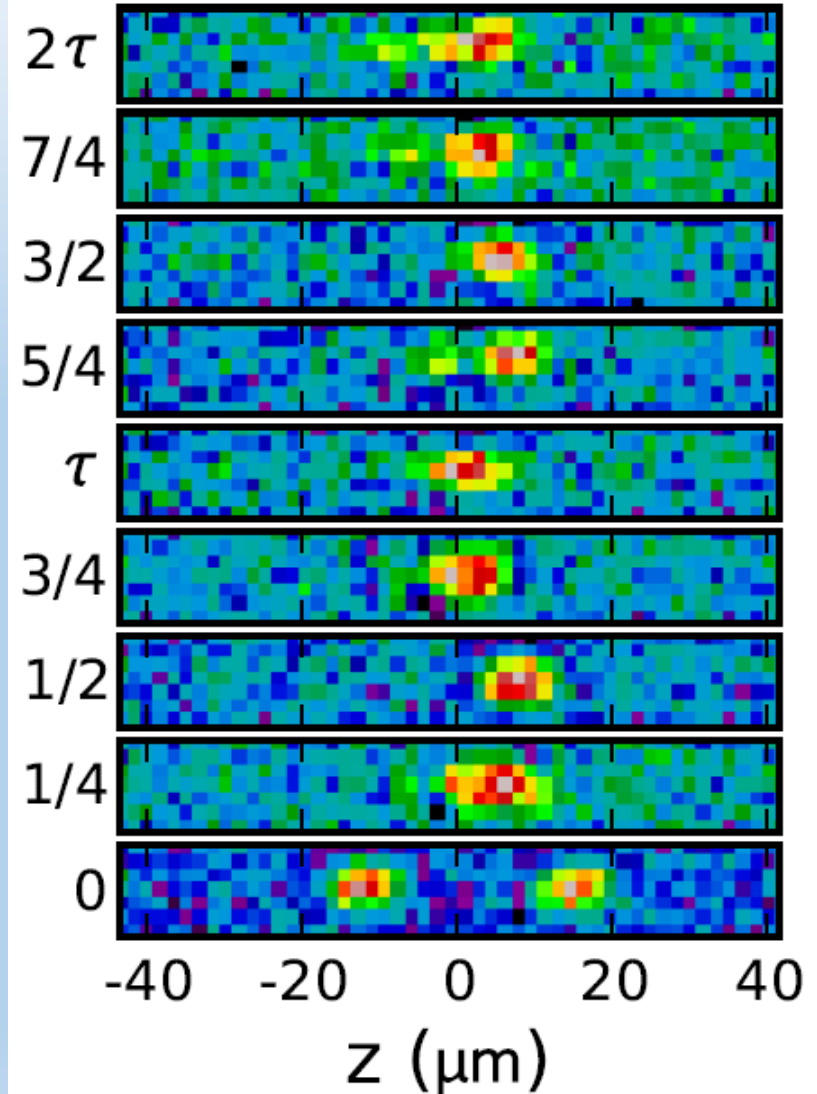
**What are the implications of this boundary?**

# On the Edge of Integrability: $N/N_c = -0.53$ , $\Delta\phi \approx 0$

Annihilation



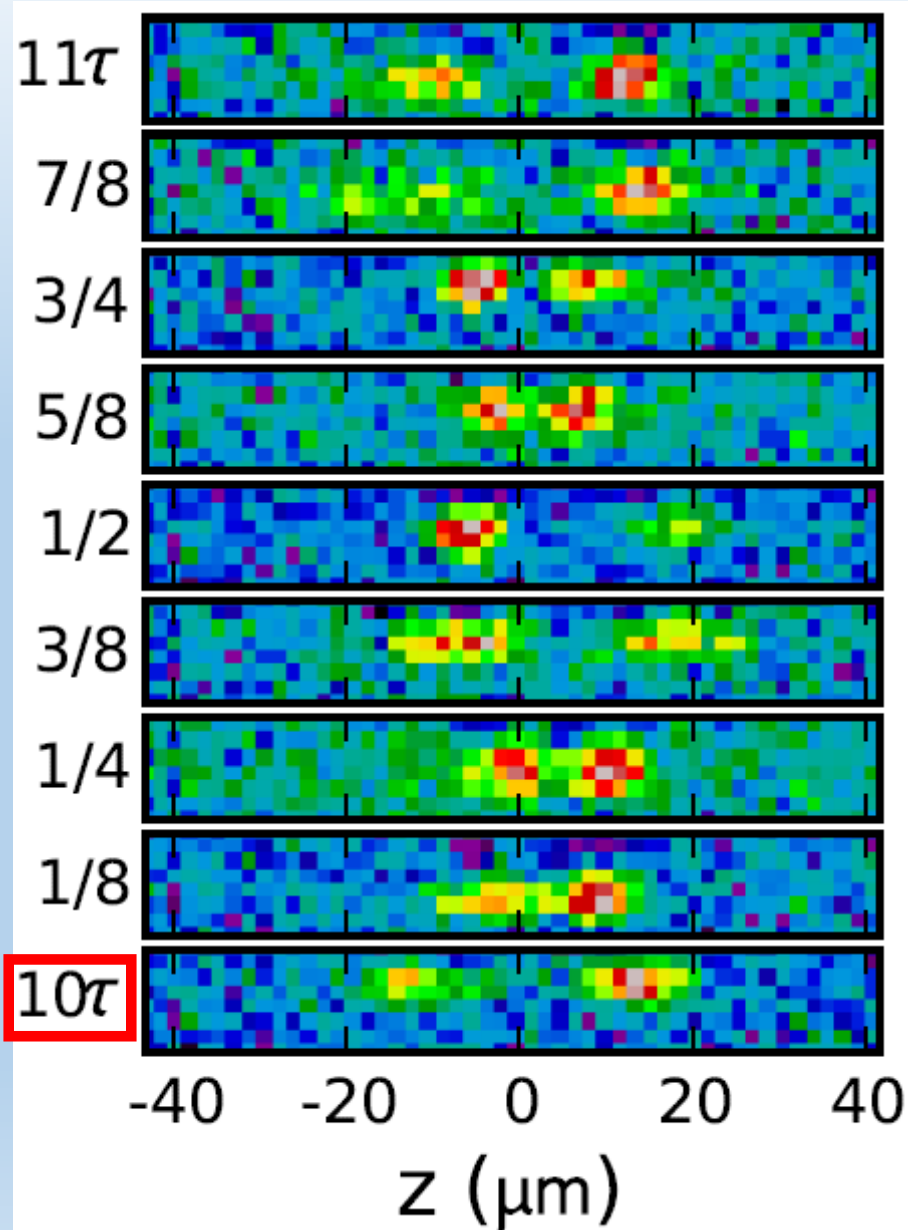
Merger



Nguyen, Dyke, Luo, Malomed & Hulet.,  
Nat. Phys. **10**, 918-922 (2014)

Simulations, Durham: Parker et al, JPB (2008)  
Billam et al, PRA (2011)

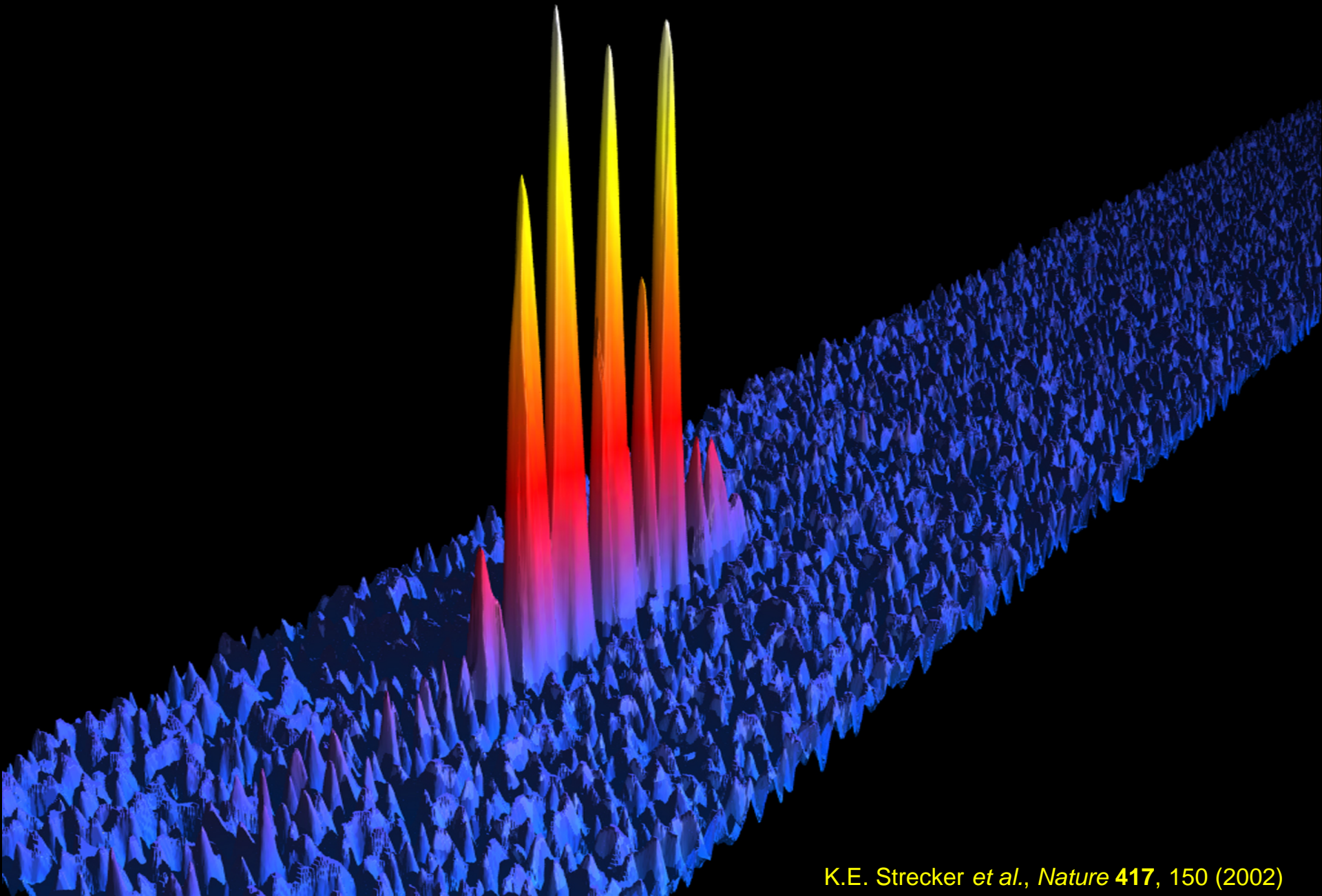
# On the Edge of Integrability: $N/N_c = -0.53$ , $\Delta\phi \approx \pi$



Behaves as if it were *integrable*:  
survives for  $> 20$  collisions



Fast quench  $\rightarrow$  Soliton Train





# Soliton Train Formed by Modulational Instability

- Soliton train formed by modulational instability (MI) following a quench:  
→ exponential growth of amplitude and phase fluctuations at

$$k_{mg} = \frac{1}{\xi}$$

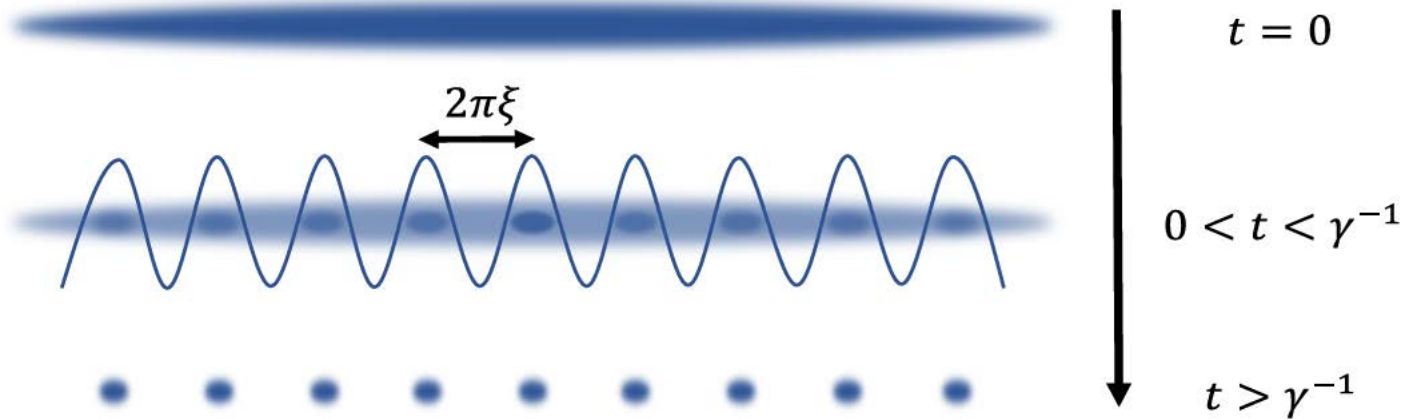
wavevector of maximal growth

$$\xi = \frac{a_r}{\sqrt{4|a_f|n_{1D}}}$$

length scale

$$\gamma = 2\omega_r|a_f|n_{1D}$$

characteristic rate



MI is seen in many wave contexts involving a self-focusing nonlinearity:

- deep water waves (Benjamin-Feir instability) – rogue waves
- plasma instabilities
- optical fibers with anomalous dispersion

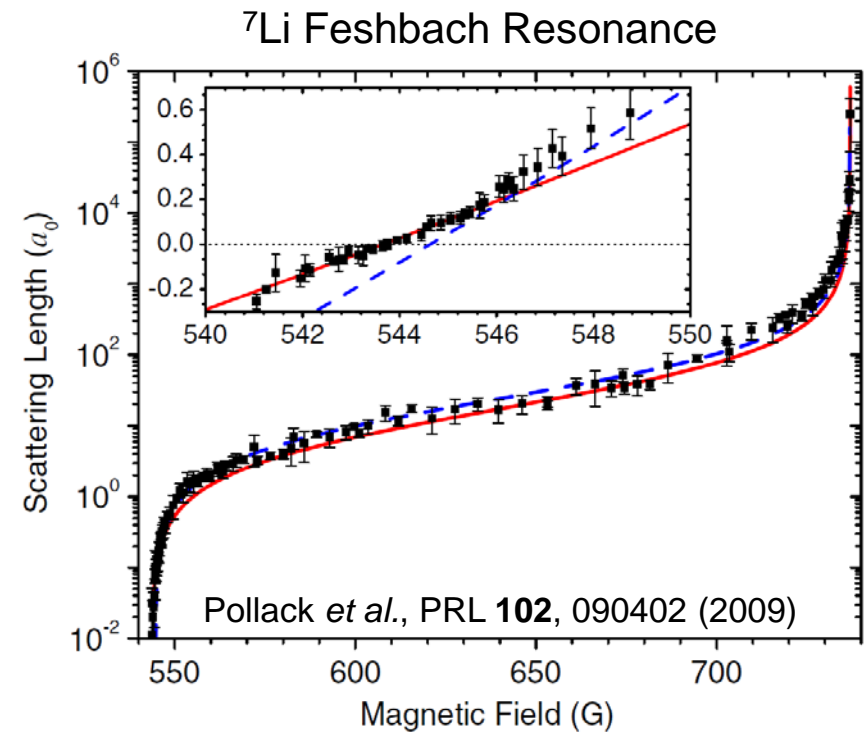
# Formation of Soliton Trains

- trap frequencies

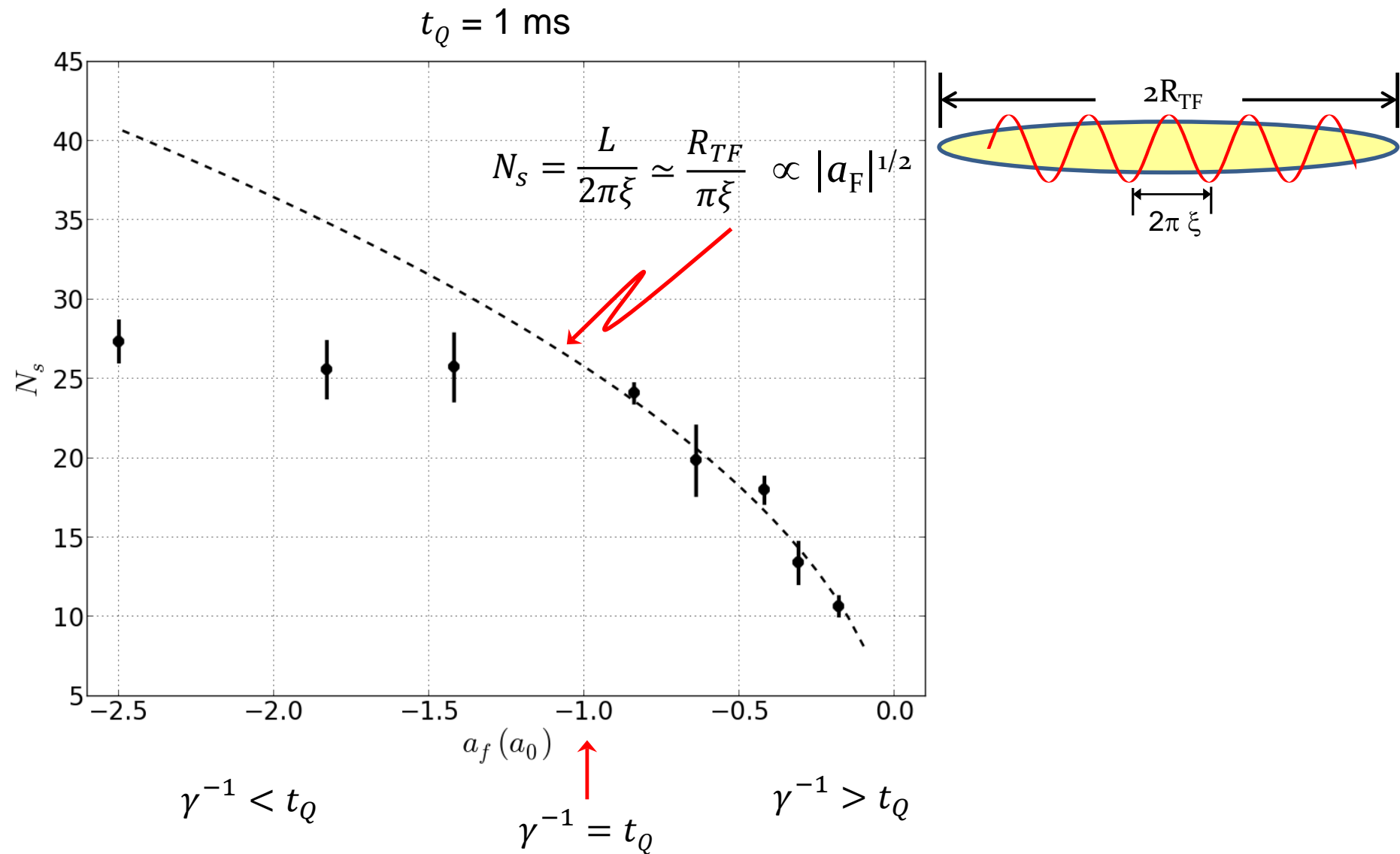
$$\omega_r = 2\pi \times 346 \text{ Hz}$$

$$\omega_z = 2\pi \times 7.4 \text{ Hz}$$

- start with BEC with  $a = +3a_0$
- quench to  $a < 0$  in  $t_Q = 1 \text{ ms}$
- hold for  $t_h$ , take *in-situ* image



# Number Solitons vs Interaction

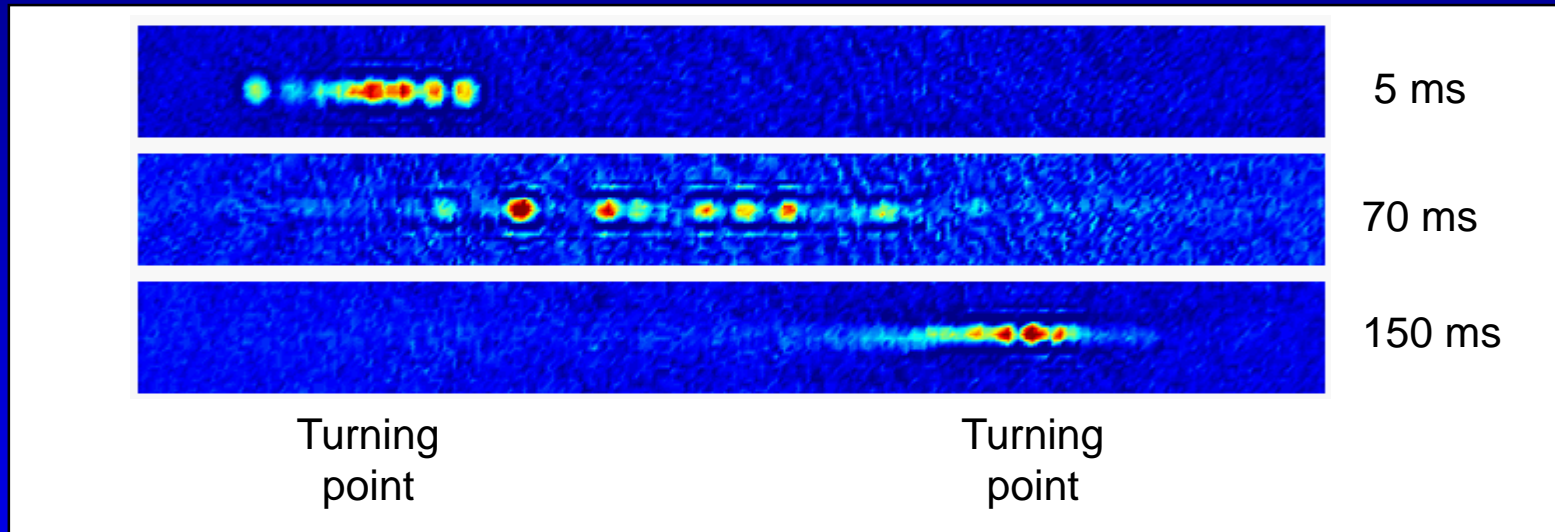


where the freeze-out time = quench time

# Solitons Appear to Repel

Create soliton train, then set into harmonic oscillation

Distance between solitons increases at bottom of well, while bunching at turning points  $\rightarrow$  repulsive interaction



K.E. Strecker et al., Nature 417, 150 (2002)

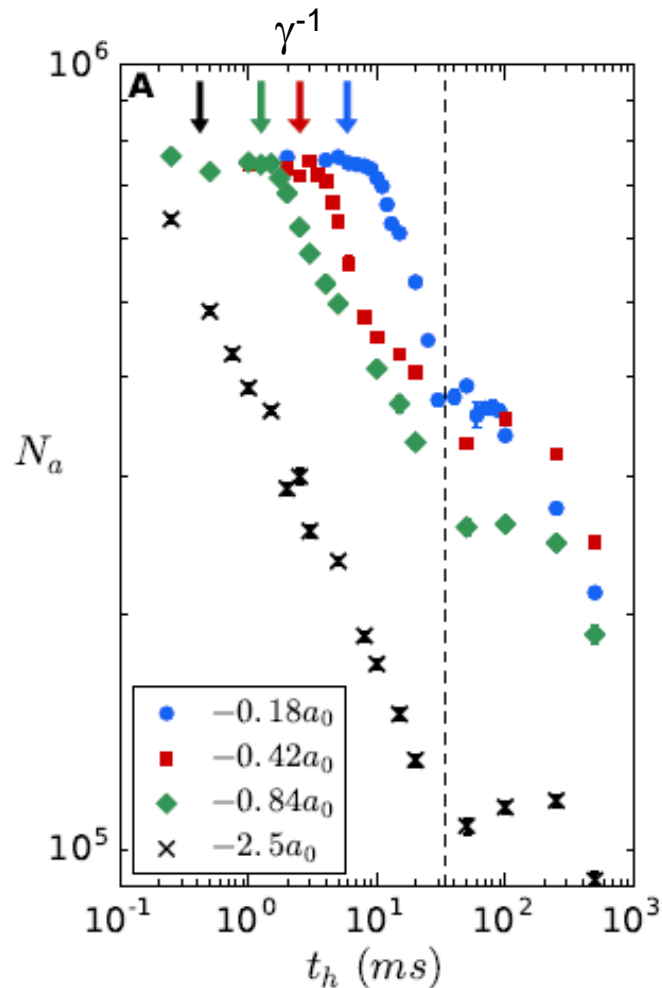
Also, S. Cornish et al., PRL 96, 170401 (2006) (JILA)

Gordon-Haus Effect: Repulsive interaction for out-of-phase solitons



How does this phase structure arise? Are they born that way, or does it develop?

# Atom Loss During Hold Time



Data collapse by plotting

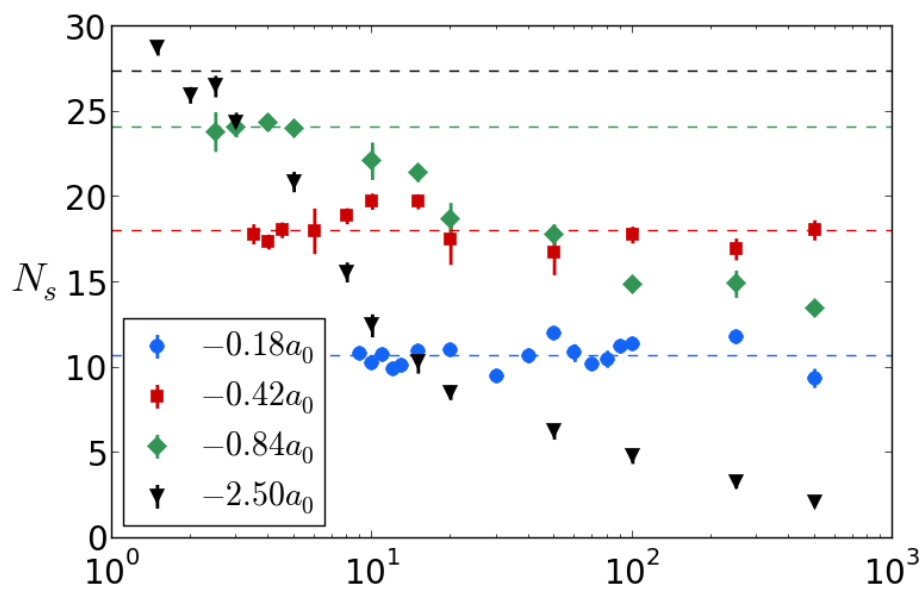
$$N_a \text{ vs. } \gamma^{-1}$$

$$N_a \propto (t_h \gamma)^\kappa$$

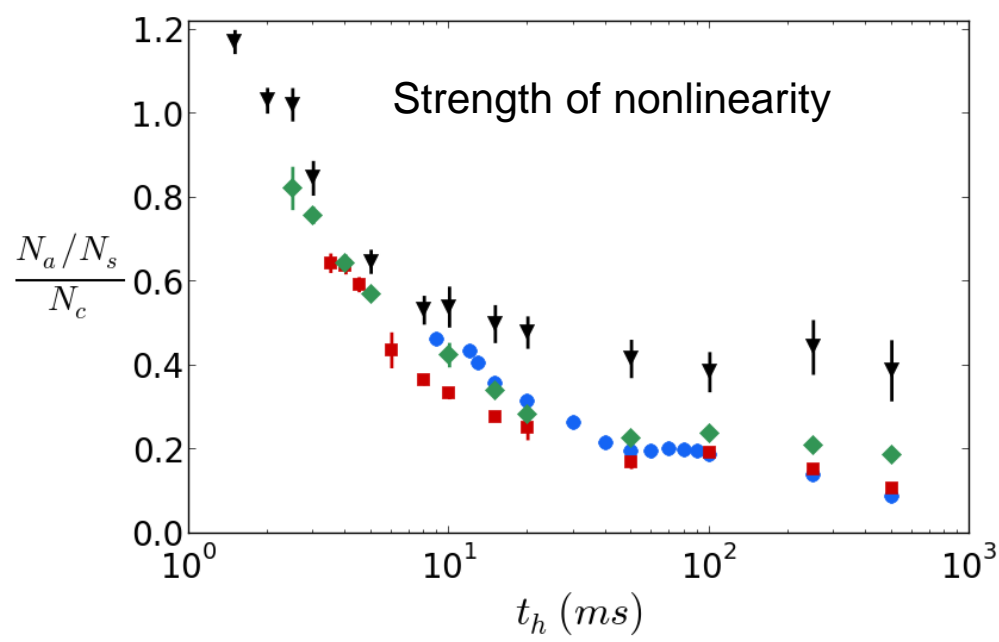
with  $\kappa = 0.35(1)$

$$t_Q = 1 \text{ ms: } \gamma^{-1} < t_Q \text{ for } a_f = -2.5a_0$$

# Soliton Number vs Hold Time



- Small  $|a_f|$ : soliton number is constant
- Large  $|a_f|$ : soliton number decreases

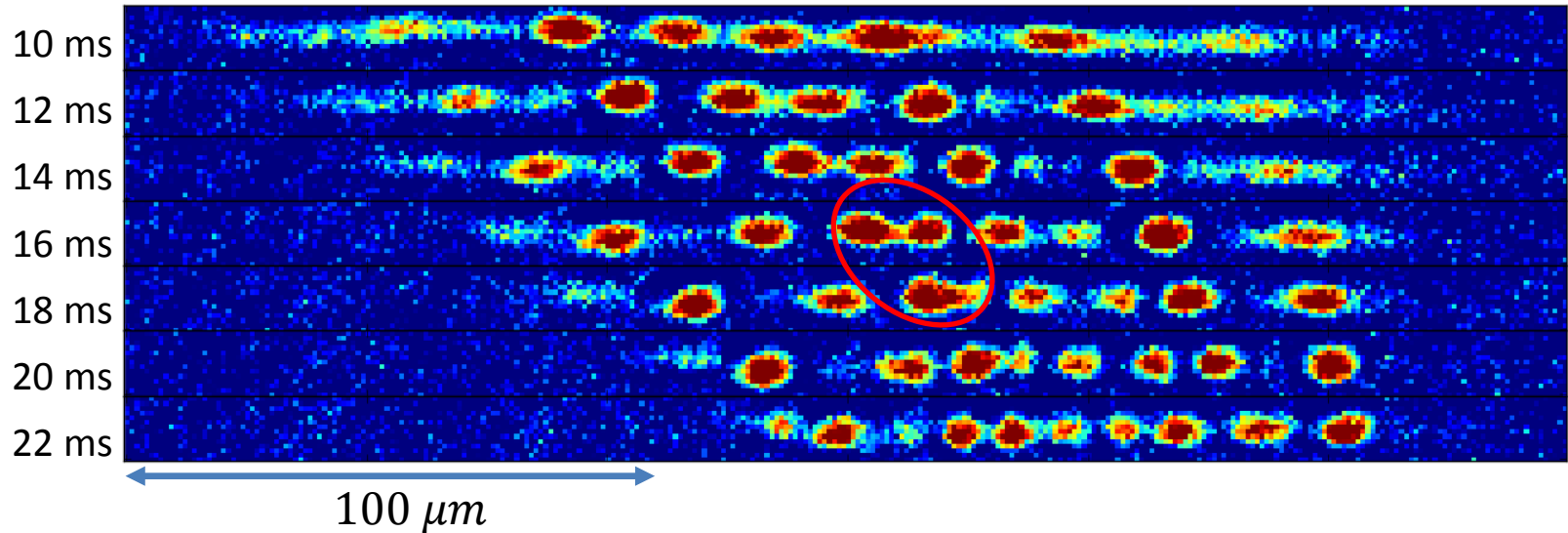


- Particle loss without loss of solitons:  
→ driven by partial, local collapse  
(Papers by: Saito/Ueda, Kagan/Shlyapnikov)

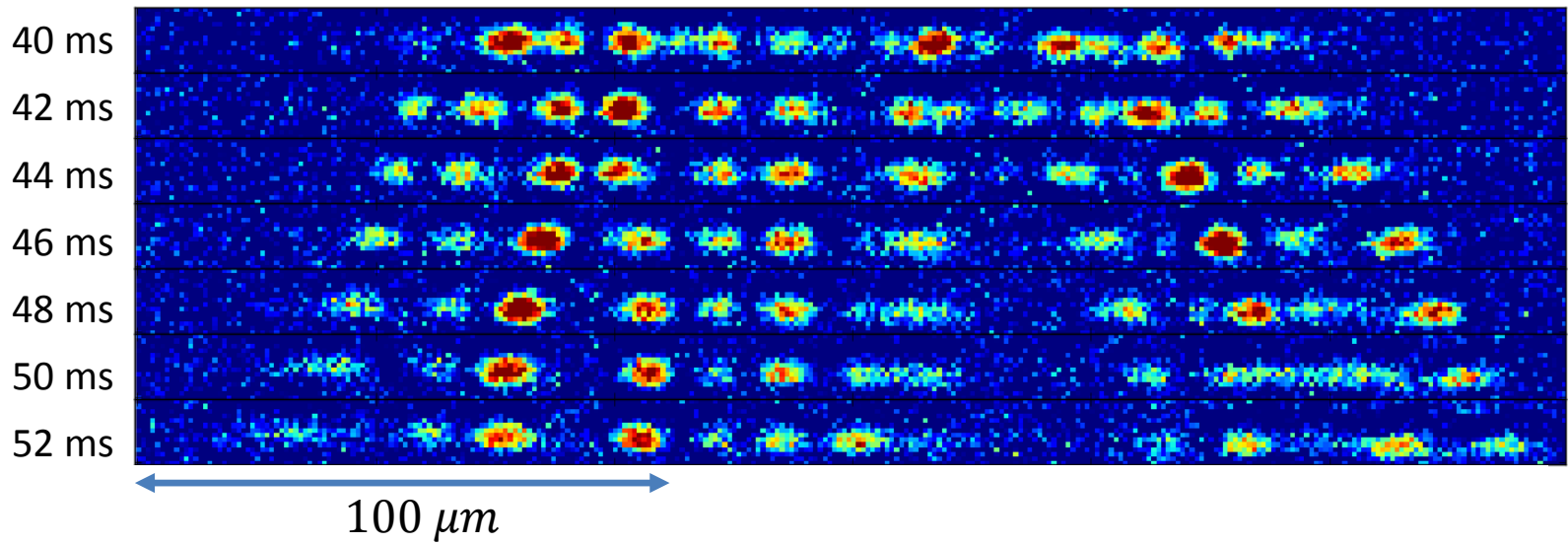
# Soliton Train Dynamics

$$a_f = -0.18a_0$$

Mainly see repulsion between solitons, with just a few cases of attraction



Breathing mode excited: compression phase at ~34 ms



# Conclusions

- System on the edge of integrability:
  - Integrability  $\rightarrow$  solitons pass through one another without changing shape, speed, or amplitude
  - Yet, they undergo phase-dependent collisions
  - Breakdown of integrability is sudden and depends on  $\Delta\phi$ : annihilation and merger
- Soliton train is “born” with an alternating phase, whereas the notion of self-assembly by soliton annihilations is not supported by observation (L. Salasnich; H. Stoof)
- Freeze-out time agrees with modulational instability (L. Carr and J. Brand)
- After the freeze-out time, there is a continuous loss of atoms driven by partial local collapses (Saito/Ueda and Kagan/Shlyapnikov)





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